

BOCHNER-WEITZENBÖCK FORMULAS AND CURVATURE ACTIONS ON RIEMANNIAN MANIFOLDS

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ABSTRACT. Gradients are natural first order differential operators depending on Riemannian metrics. The principal symbols of them are related to the enveloping algebra and higher Casimir elements. We give formulas in the enveloping algebra that induce not only identities for higher Casimir elements but also all Bochner-Weitzenböck formulas for gradients. As applications, we give some vanishing theorems.

1. INTRODUCTION

The Dirac operator is an important tool in a wide range of mathematics and physics. The principal symbols of the Dirac operator are known as the Clifford multiplications and constitute the Clifford algebra. The algebra not only gives features of the Dirac operator but also has importance by itself [19]. We can generalize the Dirac operator under the condition that operators are first order and conformally covariant [10], [22]. The generalized operators are called *gradients* or *Stein-Weiss operators* and include basic operators in Riemannian and spin geometry: the exterior derivative, the interior derivative, the conformal Killing operator, the twistor operator and the Rarita-Schwinger operator. From recent research by T. Branson et al., we know various properties and applications of gradients, which give a new direction in geometry and analysis.

Let (M, g) be an n -dimensional oriented Riemannian manifold and let $\mathbf{SO}(M)$ be the bundle of oriented orthonormal frames with structure group $\mathrm{SO}(n)$. An irreducible unitary $\mathrm{SO}(n)$ -module V_ρ with highest weight ρ gives an associated vector bundle $\mathbf{S}_\rho := \mathbf{SO}(M) \times_{\pi_\rho} V_\rho$ on M . On this vector bundle, we have a covariant derivative ∇ induced from the Levi-Civita connection. We decompose the target bundle of ∇ as

$$\mathbf{S}_\rho \otimes T_{\mathbb{C}}^*(M) = \bigoplus_{1 \leq i \leq N} \mathbf{S}_{\lambda_i}.$$

Then the derivative ∇ splits along this decomposition, and each summand is a first order differential operator,

$$D_{\lambda_i}^\rho : \Gamma(M, \mathbf{S}_\rho) \xrightarrow{\nabla} \Gamma(M, \mathbf{S}_\rho \otimes T_{\mathbb{C}}^*(M)) \xrightarrow{\text{projection}} \Gamma(M, \mathbf{S}_{\lambda_i}).$$

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We call this operator *gradient*. In [3], T. Branson discusses ellipticities of gradients by using the spectral resolution on the standard sphere. As a corollary, he shows that suitable linear combinations of gradients are bundle endomorphisms depending on the Riemannian curvature,

$$(1.1) \quad \sum_{1 \leq i \leq N} b_{\lambda_i} (D_{\lambda_i}^\rho)^* D_{\lambda_i}^\rho = \text{curvature action}.$$

These formulas are called *optimal Bochner-Weitzenböck formulas*. The vector $(b_{\lambda_1}, \dots, b_{\lambda_N})$ needs to be a solution to a certain system of linear equations (see [3, Theorem 5.10]). Since the solutions constitute an $[N/2]$ -dimensional subspace in \mathbb{R}^N , there are $[N/2]$ independent optimal Bochner-Weitzenböck formulas. This result induces vanishing theorems and eigenvalue estimates in Riemannian and spin geometry [6], [7], [8]. Furthermore, his approach of [3] is used to solve the Kato constant problem [5]. But it would be complicated to calculate the coefficients $\{b_{\lambda_i}\}_i$ and curvature actions in (1.1). On the other hand, D. Calderbank, P. Gauduchon and M. Herzlich also discuss the Kato constant problem [9]. Their key observation is that the principal symbols of gradients are related to the enveloping algebra of $\mathfrak{so}(n)$, especially higher Casimir elements. This suggests that we can reconstruct Branson's optimal Bochner-Weitzenböck formulas in terms of the enveloping algebra. In this paper, we develop the observation of [9] further, and connect higher Casimir elements to Bochner-Weitzenböck formulas. Consequently, we obtain a universal and direct construction of the coefficients $\{b_{\lambda_i}\}_i$ and curvature actions in (1.1), which makes it possible to give various vanishing theorems and eigenvalue estimates.

In Section 2, we give a short review to representations of $\mathfrak{so}(n)$. In Section 3, after we discuss the enveloping algebra and Casimir elements of $\mathfrak{so}(n)$, we give the universal Bochner-Weitzenböck formulas (Theorem 3.4). The formulas induce some identities for higher Casimir elements (Corollary 3.5). In Section 4, we discuss the principal symbols of gradients called *Clifford homomorphisms*. We relate them with the enveloping algebra by using conformal weights. From the universal Bochner-Weitzenböck formulas, we have identities for Clifford homomorphisms corresponding to (1.1) on the symbol level. We also compute eigenvalues of Casimir elements. In Section 5, we define gradients and give some fundamental properties of them. In Section 6, we introduce curvature endomorphisms corresponding to curvature actions in (1.1). An interesting observation is that the curvature endomorphism associated to the Pfaffian element depends only on the conformal Weyl tensor and the scalar curvature (Proposition 6.5). In Section 7, we give Bochner-Weitzenböck formulas for gradients (Theorem 7.1) and show that our formulas give an answer to the problem of finding out $\{b_{\lambda_i}\}_i$ in (1.1). We also discuss a relation between our Bochner-Weitzenböck formulas and Branson's conformally covariant operators in [4]. In Section 8, after mentioning gradients on spinors and differential forms, we apply our formulas to vanishing theorems for other gradients. In the last section, we discuss the four-dimensional case. Because of $\Lambda^2 = \Lambda_+^2 \oplus \Lambda_-^2$, we can calculate curvature endomorphisms more precisely and obtain vanishing theorems for higher tensor-spinor fields.

An outline of some results has been presented in a short paper [16] by the author. We discuss the details and develop them in the present paper. New results and many examples are included.

2. REPRESENTATIONS OF $\mathrm{SO}(n)$ AND $\mathrm{Spin}(n)$

In this section, we give a short review to representation theory of the special orthogonal group $\mathrm{SO}(n)$ or the spin group $\mathrm{Spin}(n)$ [18], [24]. Let \mathbb{R}^n be the n -dimensional Euclidean space with inner product $\langle \cdot, \cdot \rangle$ and standard basis $\{e_i\}_{i=1}^n$. Associating $\xi \wedge \eta$ with the skew-symmetric endomorphism

$$(\xi \wedge \eta)(u) = \langle \xi, u \rangle \eta - \langle \eta, u \rangle \xi \quad \text{for } u \in \mathbb{R}^n,$$

we identify $\Lambda^2(\mathbb{R}^n)$ with the Lie algebra $\mathfrak{so}(n)$ of $\mathrm{SO}(n)$ or $\mathrm{Spin}(n)$. Set $e_{ij} := e_i \wedge e_j$, and we know that $\{e_{ij}\}_{1 \leq i, j \leq n}$ satisfy

$$(2.1) \quad \begin{aligned} e_{ij} &= -e_{ji}, \\ [e_{kl}, e_{ij}] &= \delta_{ki}e_{lj} + \delta_{kj}e_{il} - \delta_{il}e_{kj} - \delta_{lj}e_{ik}, \end{aligned}$$

and $\{e_{ij} | 1 \leq i < j \leq n\}$ constitute a basis of $\mathfrak{so}(n)$.

We choose a Cartan subalgebra \mathfrak{h} of $\mathfrak{so}(n)$ with basis $\{e_{2i-1, 2i} | 1 \leq i \leq m = [n/2]\}$ and fix a basis $\{\mu_i\}_{1 \leq i \leq m}$ of $(\sqrt{-1}\mathfrak{h})^*$ such that $\mu_i(-\sqrt{-1}e_{2j-1, 2j}) = \delta_{ij}$. Normalizing the Killing form on $\mathfrak{so}(n)$, we have a positive definite inner product on $(\sqrt{-1}\mathfrak{h})^*$ such that $\langle \mu_i, \mu_j \rangle = \delta_{ij}$. Thus we regard $(\sqrt{-1}\mathfrak{h})^*$ as the m -dimensional Euclidean space and denote μ_i by

$$\mu_i = (\underbrace{0, \dots, 0}_{i-1}, 1, \underbrace{0, \dots, 0}_{m-i}).$$

We consider a finite-dimensional irreducible unitary representation (π, V) of $\mathrm{SO}(n)$ or $\mathrm{Spin}(n)$, and decompose the representation space V into simultaneous eigenspaces with respect to $\sqrt{-1}\mathfrak{h}$. Each eigenvalue ν called *weight* is an integral or half-integral linear combination of $\{\mu_i\}_i$, namely, $\nu = \sum \nu^i \mu_i = (\nu^1, \dots, \nu^m)$ in $\mathbb{Z}^m \cup (\mathbb{Z} + 1/2)^m$. Ordering the weights lexicographically, we have the highest weight $\rho = (\rho^1, \dots, \rho^m)$ for (π, V) with multiplicity one. This highest weight satisfies *the dominant condition*

$$\begin{aligned} \rho^1 &\geq \rho^2 \geq \dots \geq \rho^{m-1} \geq |\rho^m|, \quad \text{for } n = 2m, \\ \rho^1 &\geq \rho^2 \geq \dots \geq \rho^{m-1} \geq \rho^m \geq 0, \quad \text{for } n = 2m + 1. \end{aligned}$$

Conversely, for a dominant weight ρ in $\mathbb{Z}^m \cup (\mathbb{Z} + 1/2)^m$, we can construct a unique irreducible unitary representation with highest weight ρ up to equivalence. Therefore we denote by (π_ρ, V_ρ) an irreducible representation with highest weight ρ of $\mathrm{SO}(n)$, $\mathrm{Spin}(n)$ or $\mathfrak{so}(n)$. Note that if ρ is half-integral, then (π_ρ, V_ρ) does not factor through a representation of $\mathrm{SO}(n)$. When writing dominant weights, we denote by k_j a string of k with length j and sometimes omit a terminal string of zeros. For example, the highest weight of the spinor representation is

$$((1/2)_{m-1}, \pm 1/2) = (\underbrace{1/2, 1/2, \dots, 1/2}_{m-1}, \pm 1/2),$$

and the one of the representation on $\Lambda^p(\mathbb{R}^n) \otimes \mathbb{C}$ is

$$(1_p) = (\underbrace{1, \dots, 1}_p, \underbrace{0, \dots, 0}_{m-p}).$$

3. ENVELOPING ALGEBRA AND HIGHER CASIMIR ELEMENTS

Let $\mathfrak{so}(n, \mathbb{C})$ be the complexification of $\mathfrak{so}(n)$ and let $U(\mathfrak{so}(n, \mathbb{C}))$ be its enveloping algebra. The enveloping algebra is the quotient algebra of the tensor algebra of $\mathfrak{so}(n, \mathbb{C})$ by two-sided ideal generated by all $(X \otimes Y - Y \otimes X - [X, Y])$ for X, Y in $\mathfrak{so}(n, \mathbb{C})$. Each representation (π, V) of $\mathfrak{so}(n)$ is lifted naturally to a representation of the enveloping algebra denoted by the same notation (π, V) .

The center \mathfrak{Z} of $U(\mathfrak{so}(n, \mathbb{C}))$ is characterized as the invariant subalgebra in $U(\mathfrak{so}(n, \mathbb{C}))$ under the adjoint action of $\mathrm{SO}(n)$. We call elements in \mathfrak{Z} *Casimir elements*. From Schur's lemma, we know that every Casimir element is a constant on irreducible $\mathfrak{so}(n)$ -module.

We shall construct an algebraic basis of \mathfrak{Z} [20], [21], [24]. First, we define the usual Casimir element by $c_2 := \sum_{i,j} e_{ij}e_{ji}$. It is known that the eigenvalue of c_2 on irreducible $\mathfrak{so}(n)$ -module V_ρ is

$$(3.1) \quad \pi_\rho(c_2) = 2\langle \delta + \rho, \delta + \rho \rangle - 2\langle \delta, \delta \rangle = 2\langle \rho, \rho \rangle + 4\langle \rho, \delta \rangle,$$

where δ is half the sum of the positive roots

$$\delta = \begin{cases} (m-1, m-2, \dots, 1, 0) & \text{for } n = 2m, \\ (m-1/2, m-3/2, \dots, 3/2, 1/2) & \text{for } n = 2m+1. \end{cases}$$

For example, on the natural representation $(\pi_{\mu_1}, V_{\mu_1}) = (\pi_{\mu_1}, \mathbb{C}^n)$, we have $\pi_{\mu_1}(c_2) = 2(n-1)$.

Next we construct higher Casimir elements. For each non-negative integer q , we define an element e_{ij}^q in $U(\mathfrak{so}(n, \mathbb{C}))$ by

$$(3.2) \quad e_{ij}^q := \begin{cases} \sum_{1 \leq i_1, i_2, \dots, i_{q-1} \leq n} e_{ii_1} e_{i_1 i_2} \cdots e_{i_{q-1} j}, & q \geq 1, \\ \delta_{ij}, & q = 0. \end{cases}$$

This e_{ij}^q behaves like e_{ij} under the adjoint action of $\mathfrak{so}(n)$ as follows.

Lemma 3.1. *The elements $\{e_{ij}^q | q \in \mathbb{Z}_{\geq 0}, i, j = 1, \dots, n\}$ satisfy*

$$(3.3) \quad [e_{kl}, e_{ij}^q] = \delta_{ki} e_{lj}^q + \delta_{kj} e_{il}^q - \delta_{il} e_{kj}^q - \delta_{lj} e_{ik}^q,$$

$$(3.4) \quad \sum_{1 \leq k \leq n} e_{ik}^p e_{kj}^q = e_{ij}^{p+q}.$$

Proof. We calculate the adjoint action of $\mathfrak{so}(n)$ on e_{ij}^q . From (2.1), we have

$$\begin{aligned} [e_{kl}, e_{ij}^q] &= \sum_{i_1, \dots, i_{q-1}} [e_{kl}, e_{ii_1}] e_{i_1 i_2} \cdots e_{i_{q-1} j} + \cdots + \sum_{i_1, \dots, i_{q-1}} e_{ii_1} e_{i_1 i_2} \cdots [e_{kl}, e_{i_{q-1} j}] \\ &= \delta_{ki} e_{lj}^q + \delta_{kj} e_{il}^q - \delta_{il} e_{kj}^q - \delta_{lj} e_{ik}^q. \end{aligned}$$

The equation (3.4) is clear from the definition (3.2). \square

The equation (3.3) implies that the trace $c_q := \sum_i e_{ii}^q$ is an invariant element. Thus we have higher Casimir elements $\{c_q\}_{q \geq 0}$. Note that c_0 is equal to n and c_1 is zero. The eigenvalues of c_q on irreducible $\mathfrak{so}(n)$ -modules are calculated in Section 4.

In the case of $n = 2m+1$, Casimir elements $\{c_q\}_q$ generate the center \mathfrak{Z} algebraically. On the other hand, in the case of $n = 2m$, we need another Casimir

element to generate \mathfrak{Z} . We define pf in \mathfrak{Z} by

$$(3.5) \quad \text{pf} := \frac{1}{(\sqrt{-1})^m 2^m m!} \sum_{\sigma \in \mathfrak{S}_{2m}} \text{sign}(\sigma) e_{\sigma(1)\sigma(2)} e_{\sigma(3)\sigma(4)} \cdots e_{\sigma(2m-1)\sigma(2m)},$$

where \mathfrak{S}_{2m} is the permutation group of $\{1, \dots, 2m\}$. We call the Casimir element pf *the Pfaffian element*. The following facts on these Casimir elements are known. (See [24, pp. 368–373].)

Proposition 3.2. (1) *In the case of $n = 2m$, $\{c_2, c_4, \dots, c_{2m-2}, \text{pf}\}$ generate \mathfrak{Z} algebraically. The eigenvalue of pf on an irreducible $\mathfrak{so}(2m)$ -module V_ρ is*

$$(3.6) \quad \pi_\rho(\text{pf}) = (\rho^1 + m - 1)(\rho^2 + m - 2) \cdots (\rho^{m-1} + 1)\rho^m.$$

(2) *In the case of $n = 2m + 1$, $\{c_2, c_4, \dots, c_{2m}\}$ generate \mathfrak{Z} algebraically.*

From the above proposition, we have the problem of how to realize c_{2q+1} as a polynomial in $\{c_{2p}\}_p$. To answer it, we search to find how e_{ij}^q is related to e_{ji}^q . From (3.3), we have

$$e_{ij}^{q+1} = (1 - n)e_{ij}^q + \delta_{ji}c_q - e_{ji}^q - \sum_k e_{kj}^q e_{ki}.$$

This equation implies that e_{ij}^q is a linear combination of $\{e_{ji}^p\}_{p=0}^q$ as

$$e_{ij}^q = \sum_{p=0}^q a_{q,p} e_{ji}^p.$$

Here $\{a_{q,p}\}_{q \geq p \geq 0}$ in \mathfrak{Z} are given by a recursion formula. Since the recursion formula is complicated, we translate e_{ij}^q to another element. We define \hat{e}_{ij} by

$$\hat{e}_{ij} := e_{ij} + \frac{n-1}{2} \delta_{ij},$$

and \hat{e}_{ij}^q by

$$\hat{e}_{ij}^q := \begin{cases} \sum_{1 \leq i_1, i_2, \dots, i_{q-1} \leq n} \hat{e}_{ii_1} \hat{e}_{i_1 i_2} \cdots \hat{e}_{i_{q-1} j}, & q \geq 1, \\ \delta_{ij}, & q = 0. \end{cases}$$

Note that \hat{e}_{ij}^q is related to e_{ij}^q as

$$\hat{e}_{ij}^q = \sum_{p=0}^q \binom{q}{p} \left(\frac{n-1}{2} \right)^{q-p} e_{ij}^p.$$

We also define the translated Casimir element \hat{c}_q by $\hat{c}_q := \sum_i \hat{e}_{ii}^q$.

Lemma 3.3. *The translated elements $\{\hat{e}_{ij}^q | q \in \mathbb{Z}_{\geq 0}, i, j = 1, \dots, n\}$ satisfy*

$$(3.7) \quad [\hat{e}_{kl}, \hat{e}_{ij}^q] = \delta_{ki} \hat{e}_{lj}^q + \delta_{kj} \hat{e}_{il}^q - \delta_{il} \hat{e}_{kj}^q - \delta_{lj} \hat{e}_{ik}^q,$$

$$(3.8) \quad \sum_k \hat{e}_{ik}^p \hat{e}_{kj}^q = \hat{e}_{ij}^{p+q},$$

$$(3.9) \quad \hat{e}_{ij} = -\hat{e}_{ji} + (n-1)\delta_{ij}.$$

We particularly obtain a key relation

$$(3.10) \quad \hat{e}_{ij}^{q+1} = \delta_{ji} \hat{c}_q - \hat{e}_{ji}^q - \sum_k \hat{e}_{kj}^q \hat{e}_{ki}.$$

Proof. We can prove (3.7)–(3.9) in the same way as Lemma 3.1. It follows from (3.7)–(3.9) that

$$\begin{aligned}
\hat{e}_{ij}^{q+1} &= \sum_k ([\hat{e}_{ik}, \hat{e}_{kj}^q] + \hat{e}_{kj}^q \hat{e}_{ik}) \\
&= \sum_k \{(\delta_{ik} \hat{e}_{kj}^q + \delta_{ij} \hat{e}_{kk}^q - \delta_{kk} \hat{e}_{ij}^q - \delta_{kj} \hat{e}_{ki}^q) + \hat{e}_{kj}^q (-\hat{e}_{ki} + (n-1)\delta_{ik})\} \\
&= \hat{e}_{ij}^q + \delta_{ij} \hat{c}_q - n \hat{e}_{ij}^q - \hat{e}_{ji}^q - \sum_k \hat{e}_{kj}^q \hat{e}_{ki} + (n-1) \hat{e}_{ij}^q \\
&= \delta_{ij} \hat{c}_q - \hat{e}_{ji}^q - \sum_k \hat{e}_{kj}^q \hat{e}_{ki}.
\end{aligned}$$

Thus we obtain (3.10). \square

By using the above lemma, we connect \hat{e}_{ij}^q with \hat{e}_{ji}^q .

Theorem 3.4. *The translated element \hat{e}_{ij}^q is a linear combination of $\{\hat{e}_{ji}^p\}_{p=0}^q$ whose coefficients are Casimir elements,*

$$(3.11) \quad \hat{e}_{ij}^q = (-1)^q \hat{e}_{ji}^q - \frac{1 - (-1)^q}{2} \hat{e}_{ji}^{q-1} + \sum_{p=0}^{q-1} (-1)^p \hat{c}_{q-1-p} \hat{e}_{ji}^p.$$

Then we have

$$(3.12) \quad \hat{e}_{ij}^{2q} = \hat{e}_{ji}^{2q} + \sum_{p=0}^{2q-1} (-1)^p \hat{c}_{2q-1-p} \hat{e}_{ji}^p,$$

$$(3.13) \quad \hat{e}_{ij}^{2q+1} = -\hat{e}_{ji}^{2q+1} - \hat{e}_{ji}^{2q} + \sum_{p=0}^{2q} (-1)^p \hat{c}_{2q-p} \hat{e}_{ji}^p.$$

Proof. Setting $\hat{e}_{ij}^q = \sum_{p=0}^q \hat{a}_{q,p} \hat{e}_{ji}^p$, we shall produce a recursion formula of $\{\hat{a}_{q,p}\}_{q \geq p \geq 0}$ in §. It follows from (3.10) that

$$\begin{aligned}
\hat{e}_{ij}^{q+1} &= \delta_{ji} \hat{c}_q - \hat{e}_{ji}^q - \sum_k \hat{e}_{kj}^q \hat{e}_{ki} \\
&= \delta_{ji} \hat{c}_q - \hat{e}_{ji}^q - \sum_k \sum_p \hat{a}_{q,p} \hat{e}_{jk}^p \hat{e}_{ki} \\
&= \delta_{ji} \hat{c}_q - \hat{e}_{ji}^q - \sum_{p=0}^q \hat{a}_{q,p} \hat{e}_{ji}^{p+1} \\
&= -\hat{a}_{q,q} \hat{e}_{ji}^{q+1} + (-\hat{a}_{q,q-1} - 1) \hat{e}_{ji}^q - \sum_{p=0}^{q-2} \hat{a}_{q,p} \hat{e}_{ji}^{p+1} + \hat{c}_q \delta_{ji} \\
&= \sum_{p=0}^{q+1} \hat{a}_{q+1,p} \hat{e}_{ji}^p.
\end{aligned}$$

Then we have a recursion formula for $\{\hat{a}_{q,p}\}_{q \geq p \geq 0}$,

$$\hat{a}_{q+1,p} = \begin{cases} -\hat{a}_{q,q}, & p = q+1, \\ -\hat{a}_{q,q-1} - 1, & p = q, \\ -\hat{a}_{q,p-1}, & 1 \leq p \leq q-1, \\ \hat{c}_q, & p = 0. \end{cases}$$

Because of $\hat{e}_{ij}^0 = \hat{e}_{ji}^0$ and $\hat{e}_{ij}^1 = -\hat{e}_{ji}^1 + (n-1)\hat{e}_{ij}^0$, the initial condition of $\hat{a}_{q,p}$ is $(\hat{a}_{0,0}, \hat{a}_{1,0}, \hat{a}_{1,1}) = (1, n-1, -1)$. Then

$$\hat{a}_{q+1,p} = \begin{cases} (-1)^{q+1}, & p = q+1, \\ (-1)^q(n-1) - \frac{1-(-1)^q}{2} = (-1)^q\hat{c}_0 - \frac{1-(-1)^{q+1}}{2}, & p = q, \\ (-1)^p\hat{c}_{q-p}, & 0 \leq p \leq q-1. \end{cases}$$

Thus we have proved the theorem. \square

In Section 6, we show that (3.11) induces Bochner-Weitzenböck formulas on Riemannian manifolds. Hence we call (3.11) *the universal Bochner-Weitzenböck formula*.

Take the trace in (3.13), and we have identities for $\{\hat{c}_q\}_{q \geq 0}$.

Corollary 3.5. *The Casimir elements $\{\hat{c}_0, \hat{c}_1, \dots\}$ satisfy*

$$2\hat{c}_{2q+1} = -\hat{c}_{2q} + \sum_{p=0}^{2q} (-1)^p \hat{c}_{2q-p} \hat{c}_p$$

for $q = 0, 1, \dots$.

D. Calderbank pointed out to the author how the formula (3.11) was related with [9]. We set

$$(3.14) \quad \hat{E}_{ij}^q := -\frac{1+(-1)^q}{2} \hat{e}_{ij}^q + \sum_{p=0}^q (-1)^p \hat{c}_{q-p} \hat{e}_{ij}^p.$$

Then the universal Bochner-Weitzenböck formula means $\hat{E}_{ij}^q = (-1)^q \hat{E}_{ji}^q$. This symmetry is the same as that of a formula in [9, Theorem 4.8]. A better point of our formula is that (3.11) is independent of representations. We note that a formula with such symmetry in the enveloping algebra implies a Bochner-Weitzenböck formula. We actually have another formula related to the Pfaffian element independent of (3.11) in the next section.

4. PRINCIPAL SYMBOLS OF GRADIENTS

We consider the Clifford multiplication on the spinor space V_Δ , where n is odd and Δ is $((1/2)_m)$. The Clifford multiplication is an action of ξ in \mathbb{R}^n on V_Δ ,

$$V_\Delta \ni \phi \mapsto \xi \cdot \phi \in V_\Delta,$$

and satisfies the Clifford relation $\xi \cdot \eta + \eta \cdot \xi = -2\langle \xi, \eta \rangle \text{id}$. To generalize the Clifford multiplication, we use another definition as follows. We consider the tensor representation $(\pi_\Delta \otimes \pi_{\mu_1}, V_\Delta \otimes (\mathbb{R}^n \otimes \mathbb{C}))$ and decompose it into irreducible modules, $V_\Delta \otimes \mathbb{C}^n = V_T \oplus V_\Delta$. Here T is $(3/2, (1/2)_{m-1})$. We denote by Π_Δ^Δ the orthogonal projection from $V_\Delta \otimes \mathbb{C}^n$ onto V_Δ and show that $\Pi_\Delta^\Delta(\phi \otimes \xi)$ is a constant multiple of $\xi \cdot \phi$. Thus the projection Π_Δ^Δ gives another definition of the Clifford multiplication.

We shall generalize the above discussion to other representation spaces. We consider the tensor representation $(\pi_\rho \otimes \pi_{\mu_1}, V_\rho \otimes \mathbb{C}^n)$ and its irreducible decomposition

$$V_\rho \otimes \mathbb{C}^n = \bigoplus_{\lambda} V_{\lambda}.$$

The highest weights of irreducible components occur with multiplicity one and are characterized as follows.

Proposition 4.1 ([10, Theorem 3.4]).

- (1) When $n = 2m$, or when $n = 2m + 1$ and $\rho^m = 0$, the highest weight of an irreducible component in $V_\rho \otimes \mathbb{C}^n$ is dominant, and $\rho \pm \mu_i$ for $i = 1, 2, \dots, m$.
- (2) When $n = 2m + 1$ and $\rho^m > 0$, the highest weight is dominant, and ρ or $\rho \pm \mu_i$ for $i = 1, 2, \dots, m$.

The tensor inner product on $V_\rho \otimes \mathbb{C}^n$ induces the one on V_λ . Hence each component is orthogonal to others. We denote by Π_λ^ρ the orthogonal projection from $V_\rho \otimes \mathbb{C}^n$ onto V_λ .

Definition 4.2. For ξ in \mathbb{C}^n , we define a linear mapping $p_\lambda^\rho(\xi)$ from V_ρ to V_λ by

$$\mathbb{C}^n \times V_\rho \ni (\xi, \phi) \mapsto p_\lambda^\rho(\xi)\phi := \Pi_\lambda^\rho(\phi \otimes \xi) \in V_\lambda.$$

We denote by $p_\lambda^\rho(\xi)^*$ the adjoint operator of $p_\lambda^\rho(\xi)$ with respect to the inner products on V_ρ and V_λ . We call these linear mappings $p_\lambda^\rho(\xi)$ and $p_\lambda^\rho(\xi)^*$ the Clifford homomorphisms associated to ρ and λ .

We shall investigate properties of Clifford homomorphisms.

Lemma 4.3. The Clifford homomorphism p_λ^ρ satisfies

$$(4.1) \quad \sum_i p_\lambda^\rho(e_i) \pi_\rho(e_{ij}) = w(\rho; \lambda) p_\lambda^\rho(e_j)$$

for each j . Here, $w(\rho; \lambda)$ is given by

$$w(\rho; \lambda) := \frac{1}{2}(\langle \delta + \lambda, \delta + \lambda \rangle - \langle \delta + \rho, \delta + \rho \rangle - n + 1).$$

We call this constant $w(\rho; \lambda)$ the conformal weight associated to ρ and λ .

Proof. We define an operator C on $V_\rho \otimes \mathbb{C}^n$ by

$$C := \pi_\rho \otimes \pi_{\mu_1}(c_2) - \pi_\rho(c_2) \otimes \text{id} - \text{id} \otimes \pi_{\mu_1}(c_2).$$

Since $\pi_\rho \otimes \pi_{\mu_1}(c_2)$ is $\pi_\lambda(c_2)$ on an irreducible component V_λ , we show from (3.1) that C is $4w(\rho; \lambda)\text{id}$ on V_λ . Then we have

$$C(\phi \otimes e_i) = C \left(\sum_{\lambda} p_\lambda^\rho(e_i) \phi \right) = \sum_{\lambda} 4w(\rho; \lambda) p_\lambda^\rho(e_i) \phi$$

for $\phi \otimes e_i$ in $V_\rho \otimes \mathbb{C}^n$. On the other hand, we get

$$\begin{aligned} C &= \pi_\rho \otimes \pi_{\mu_1}(c_2) - \pi_\rho(c_2) \otimes \text{id} - \text{id} \otimes \pi_{\mu_1}(c_2) \\ &= \sum_{ij} (\pi_\rho(e_{ij}) \otimes \text{id} + \text{id} \otimes \pi_{\mu_1}(e_{ij})) (\pi_\rho(e_{ji}) \otimes \text{id} + \text{id} \otimes \pi_{\mu_1}(e_{ji})) \\ &\quad - \sum_{ij} \pi_\rho(e_{ij}) \pi_\rho(e_{ji}) \otimes \text{id} - \sum_{ij} \text{id} \otimes \pi_{\mu_1}(e_{ij}) \pi_{\mu_1}(e_{ji}) \\ &= 2 \sum_{ij} \pi_\rho(e_{ij}) \otimes \pi_{\mu_1}(e_{ji}). \end{aligned}$$

Then,

$$\begin{aligned} C(\phi \otimes e_i) &= 2 \sum_{kl} \pi_\rho(e_{kl}) \phi \otimes \pi_{\mu_1}(e_{lk}) e_i = 2 \sum_{kl} \pi_\rho(e_{kl}) \phi \otimes (\delta_{il} e_k - \delta_{ki} e_l) \\ &= 4 \sum_k \pi_\rho(e_{ki}) \phi \otimes e_k = 4 \sum_\lambda \sum_k p_\lambda^\rho(e_k) \pi_\rho(e_{ki}) \phi. \end{aligned}$$

Therefore we have $\sum_k p_\lambda^\rho(e_k) \pi_\rho(e_{ki}) = w(\rho; \lambda) p_\lambda^\rho(e_i)$ for each λ . \square

From this lemma, we can relate Clifford homomorphisms to the enveloping algebra. The following proposition has been known in [9] and [21].

Proposition 4.4. *The Clifford homomorphisms $\{p_\lambda^\rho\}_\lambda$ satisfy*

$$(4.2) \quad \sum_\lambda w(\rho; \lambda)^q p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_j) = \pi_\rho(e_{ij}^q)$$

for $q = 0, 1, \dots$, and $i, j = 1, \dots, n$. In particular, we have

$$\sum_\lambda w(\rho; \lambda)^q \sum_i p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_i) = \pi_\rho(c_q).$$

Proof. For ϕ and ψ in V_ρ , we have

$$\delta_{ij} \langle \phi, \psi \rangle = \langle \phi \otimes e_i, \psi \otimes e_j \rangle = \sum_\lambda \langle p_\lambda^\rho(e_i) \phi, p_\lambda^\rho(e_j) \psi \rangle = \left\langle \sum_\lambda p_\lambda^\rho(e_j)^* p_\lambda^\rho(e_i) \phi, \psi \right\rangle.$$

Then we have proved (4.2) for $q = 0$,

$$(4.3) \quad \sum_\lambda p_\lambda^\rho(e_j)^* p_\lambda^\rho(e_i) = \delta_{ji}.$$

By using this equation and (4.1), we can prove (4.2) inductively. \square

We calculate the conformal weights for $\lambda = \rho, \rho \pm \mu_i$,

$$(4.4) \quad \begin{cases} w(\rho; \rho + \mu_i) = \rho^i + 1 - i, & i = 1, \dots, m, \\ w(\rho; \rho - \mu_i) = -\rho^i - n + i + 1, & i = 1, \dots, m, \\ w(\rho; \rho) = -(n-1)/2. \end{cases}$$

We assume that the number of irreducible components is N , that is,

$$N = \#\{\lambda | V_\lambda \subset V_\rho \otimes \mathbb{C}^n\}.$$

Arranging them lexicographically as $\lambda_1 = \rho + \mu_1 > \lambda_2 > \dots > \lambda_N$, we show from (4.4) that

$$w(\rho; \lambda_1) > w(\rho; \lambda_2) > \dots > w(\rho; \lambda_N)$$

except for the following case. When $n = 2m$ and $\rho^{m-1} > \rho^m = 0$, there always exist the highest weights $\lambda_+ := \rho + \mu_m$ and $\lambda_- := \rho - \mu_m$ whose conformal weights coincide. Then we have

$$w(\rho; \lambda_1) > w(\rho; \lambda_2) > \dots > w(\rho; \lambda_+) = w(\rho; \lambda_-) > \dots > w(\rho; \lambda_N).$$

We call this case *the exceptional case*. Thus the conformal weights almost differ from each other. It follows from Proposition 4.4 that we obtain

$$(p_{\lambda_1}^\rho(e_i)^* p_{\lambda_1}^\rho(e_j), \dots, p_{\lambda_N}^\rho(e_i)^* p_{\lambda_N}^\rho(e_j)) W^t = (\delta_{ij}, \pi_\rho(e_{ij}), \dots, \pi_\rho(e_{ij}^{N-1})),$$

where W is an $N \times N$ Vandermonde matrix

$$\begin{pmatrix} 1 & 1 & \cdots & 1 \\ w(\rho; \lambda_1) & w(\rho; \lambda_2) & \cdots & w(\rho; \lambda_N) \\ \cdots & \cdots & \cdots & \cdots \\ w(\rho; \lambda_1)^{N-1} & w(\rho; \lambda_2)^{N-1} & \cdots & w(\rho; \lambda_N)^{N-1} \end{pmatrix}.$$

Since W is invertible unless it is the exceptional case, the Clifford homomorphism $p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_j)$ is a linear combination of $\{\pi_\rho(e_{ij}^q)\}_q$.

Remark 4.5. It was first realized in [3, pp. 366–368] that the Vandermonde systems of conformal weights were crucial. Moreover, some formulas in [9] are based on the systems.

For the exceptional case, we need the Pfaffian element pf in (3.5) to distinguish $p_{\lambda_+}^\rho$ from $p_{\lambda_-}^\rho$.

Definition 4.6. We define an element pf_{ij} in the enveloping algebra for $i, j = 1, \dots, 2m$ by

$$\text{pf}_{ij} := \begin{cases} \text{pf}, & i = j, \\ (-1)^{i+j} \frac{2m}{(\sqrt{-1})^m 2^m m!} \sum_{\sigma \in \mathfrak{S}_{2m}^{ij}} \text{sgn}(\sigma) e_{\sigma(1)\sigma(2)} \cdots e_{\sigma(2m-1)\sigma(2m)}, & i < j, \\ -\text{pf}_{ji}, & i > j, \end{cases}$$

where \mathfrak{S}_{2m}^{ij} is the permutation group of $\{1, \dots, 2m\} \setminus \{i, j\}$.

Example 4.7 (four-dimensional case). Let $\{e_i\}_{i=1}^4$ be an oriented orthonormal basis of \mathbb{R}^4 . Then we have

$$(4.5) \quad \begin{aligned} \text{pf}_{12} &= e_{34}, & \text{pf}_{13} &= -e_{24}, & \text{pf}_{14} &= e_{23}, \\ \text{pf}_{23} &= e_{14}, & \text{pf}_{24} &= -e_{13}, & \text{pf}_{34} &= e_{12}. \end{aligned}$$

In other words, pf_{ij} is equal to $*e_{ij}$ for $i \neq j$, where we denote by $*$ the Hodge operator.

From the definition of pf_{ij} , we have

Proposition 4.8. *The elements $\{\text{pf}_{ij}\}_{i,j}$ satisfy*

$$(4.6) \quad \text{pf}_{ij} + \text{pf}_{ji} = 2\delta_{ij}\text{pf}$$

and $\sum_i \text{pf}_{ii} = 2m\text{pf}$.

The equation (4.6) gives the anti-symmetric identity $\text{pf}_{ij} - \delta_{ij}\text{pf} = -(\text{pf}_{ji} - \delta_{ji}\text{pf})$, the same as (3.14). This induces Bochner-Weitzenböck formulas in Section 6. Hence we call (4.6) the universal Bochner-Weitzenböck formula as well as (3.11).

Proposition 4.9. *We can connect pf_{ij} to the Clifford homomorphisms $\{p_\lambda^\rho\}_\lambda$,*

$$(4.7) \quad \sum_\lambda \pi_\lambda(\text{pf}) p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_j) = \pi_\rho(\text{pf}_{ij}).$$

In particular, we have

$$(4.8) \quad \sum_\lambda \pi_\lambda(\text{pf}) \sum_i p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_i) = 2m\pi_\rho(\text{pf}).$$

Proof. For ϕ and ψ in V_ρ ,

$$(4.9) \quad \begin{aligned} \langle \pi_\rho \otimes \pi_{\mu_1}(\text{pf})(\phi \otimes e_j), \psi \otimes e_i \rangle &= \sum_\lambda \langle \pi_\lambda(\text{pf})p_\lambda^\rho(e_j)\phi, p_\lambda^\rho(e_i)\psi \rangle \\ &= \left\langle \sum_\lambda \pi_\lambda(\text{pf})p_\lambda^\rho(e_i)^*p_\lambda^\rho(e_j)\phi, \psi \right\rangle. \end{aligned}$$

On the other hand, we know

$$\begin{aligned} &\pi_\rho \otimes \pi_{\mu_1}(e_{\sigma(2l-1)\sigma(2l)})(\phi \otimes e_j) \\ &= \begin{cases} \pi_\rho(e_{\sigma(2l-1)\sigma(2l)})\phi \otimes e_j + \phi \otimes e_{\sigma(2l)}, & \sigma(2l-1) = j, \\ \pi_\rho(e_{\sigma(2l-1)\sigma(2l)})\phi \otimes e_j - \phi \otimes e_{\sigma(2l-1)}, & \sigma(2l) = j, \\ \pi_\rho(e_{\sigma(2l-1)\sigma(2l)})\phi \otimes e_j, & \text{otherwise.} \end{cases} \end{aligned}$$

Then taking account of $\langle e_k, e_l \rangle = \delta_{kl}$, we have

$$\begin{aligned} &(\sqrt{-1})^m 2^m m! \langle \pi_\rho \otimes \pi_{\mu_1}(\text{pf})(\phi \otimes e_j), \psi \otimes e_i \rangle \\ &= \left\langle \pi_\rho \otimes \pi_{\mu_1} \left(\sum_{\sigma \in \mathfrak{S}_{2m}} \text{sgn}(\sigma) e_{\sigma(1)\sigma(2)} \cdots e_{\sigma(2m-1)\sigma(2m)} \right) (\phi \otimes e_j), \psi \otimes e_i \right\rangle \\ &= \begin{cases} (\sqrt{-1})^m 2^m m! \langle \pi_\rho(\text{pf})\phi, \psi \rangle, & i = j, \\ (-1)^{i+j} 2m \left\langle \pi_\rho \left(\sum_{\sigma \in \mathfrak{S}_{2m}^{ij}} \text{sgn}(\sigma) e_{\sigma(1)\sigma(2)} \cdots e_{\sigma(2m-1)\sigma(2m)} \right) \phi, \psi \right\rangle, & i < j, \\ (-1)^{i+j+1} 2m \left\langle \pi_\rho \left(\sum_{\sigma \in \mathfrak{S}_{2m}^{ij}} \text{sgn}(\sigma) e_{\sigma(1)\sigma(2)} \cdots e_{\sigma(2m-1)\sigma(2m)} \right) \phi, \psi \right\rangle, & i > j. \end{cases} \end{aligned}$$

Combining this equation and (4.9), we have (4.7). The equation (4.8) is clear from (4.6). \square

We consider the exceptional case. It follows from (3.6) that $\pi_{\lambda_+}(\text{pf}) = -\pi_{\lambda_-}(\text{pf}) \neq 0$ and $\pi_\lambda(\text{pf}) = 0$ for $\lambda \neq \lambda_\pm$. Then

$$p_{\lambda_+}^\rho(e_i)^*p_{\lambda_+}^\rho(e_j) - p_{\lambda_-}^\rho(e_i)^*p_{\lambda_-}^\rho(e_j) = \frac{1}{\pi_{\lambda_+}(\text{pf})}\pi_\rho(\text{pf}_{ij}).$$

As a result, we have the following corollary of Propositions 4.4 and 4.9.

Corollary 4.10. *We can realize $p_\lambda^\rho(e_i)^*p_\lambda^\rho(e_j)$ as a linear combination of $\{\pi_\rho(e_{ij}^q)\}_q$ and $\pi_\rho(\text{pf}_{ij})$.*

This corollary implies that we may investigate e_{ij}^q and pf_{ij} instead of Clifford homomorphisms. In fact, the universal Bochner-Weitzenböck formulas (3.11) and (4.6) give a lot of identities for Clifford homomorphisms in the following way.

We consider the equation (4.2) for $q = 0, 1$. Because of $e_{ij} = -e_{ji}$, we have

$$(4.10) \quad \sum_\lambda (p_\lambda^\rho(e_i)^*p_\lambda^\rho(e_j) + p_\lambda^\rho(e_j)^*p_\lambda^\rho(e_i)) = 2\delta_{ij},$$

$$(4.11) \quad \sum_\lambda w(\rho; \lambda)(p_\lambda^\rho(e_i)^*p_\lambda^\rho(e_j) + p_\lambda^\rho(e_j)^*p_\lambda^\rho(e_i)) = 0.$$

Note that the Clifford relation on spinor space follows from the above two equations. To construct further identities, we use (3.11) and (4.6). We define the translated

conformal weight $\hat{w}(\rho; \lambda)$ by

$$\hat{w}(\rho; \lambda) := w(\rho; \lambda) + \frac{n-1}{2}.$$

Then it is simple to show that

$$\sum_{\lambda} \hat{w}(\rho; \lambda)^q p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(e_j) = \pi_{\rho}(\hat{e}_{ij}^q).$$

Substituting (3.12) for this equation, we have

$$\begin{aligned} & \sum_{\lambda} \hat{w}(\rho; \lambda)^{2q} p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(e_j) \\ &= \sum_{\lambda} \left\{ \hat{w}(\rho; \lambda)^{2q} + \sum_{p=0}^{2q-1} (-\hat{w}(\rho; \lambda))^p \pi_{\rho}(\hat{e}_{2q-1-p}) \right\} p_{\lambda}^{\rho}(e_j)^* p_{\lambda}^{\rho}(e_i), \end{aligned}$$

and hence,

$$\sum_{\lambda} \left\{ \sum_{p=0}^{2q-1} (-\hat{w}(\rho; \lambda))^p \pi_{\rho}(\hat{e}_{2q-1-p}) \right\} (p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(e_j) + p_{\lambda}^{\rho}(e_j)^* p_{\lambda}^{\rho}(e_i)) = 0.$$

From (4.6) and (4.7) for $n = 2m$, we also have

$$\sum_{\lambda} \pi_{\lambda}(\text{pf}) (p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(e_j) + p_{\lambda}^{\rho}(e_j)^* p_{\lambda}^{\rho}(e_i)) = 2\pi_{\rho}(\text{pf})\delta_{ij}.$$

We can easily show from (3.6) and (4.4) that

$$(4.12) \quad (w(\rho; \lambda) + m - 1)\pi_{\lambda}(\text{pf}) = (w(\rho; \lambda) + m)\pi_{\rho}(\text{pf})$$

for each λ . Then

$$\begin{aligned} & \sum_{\lambda} \pi_{\lambda}(\text{pf}) w(\rho; \lambda) (p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(e_j) + p_{\lambda}^{\rho}(e_j)^* p_{\lambda}^{\rho}(e_i)) \\ &= \sum_{\lambda} (-(m-1)\pi_{\lambda}(\text{pf}) + w(\rho; \lambda)\pi_{\rho}(\text{pf}) + m\pi_{\rho}(\text{pf})) (p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(e_j) + p_{\lambda}^{\rho}(e_j)^* p_{\lambda}^{\rho}(e_i)) \\ &= -2(m-1)\pi_{\rho}(\text{pf})\delta_{ij} + 2m\pi_{\rho}(\text{pf})\delta_{ij} = 2\pi_{\rho}(\text{pf})\delta_{ij}. \end{aligned}$$

As a result, we have

$$\sum_{\lambda} \pi_{\lambda}(\text{pf}) (w(\rho; \lambda) - 1) (p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(e_j) + p_{\lambda}^{\rho}(e_j)^* p_{\lambda}^{\rho}(e_i)) = 0.$$

Thus we obtain algebraic relations among Clifford homomorphisms like the Clifford relation.

Theorem 4.11. *Let ξ and η be in \mathbb{R}^n . Then the Clifford homomorphisms $\{p_{\lambda}^{\rho}\}_{\lambda}$ satisfy*

$$\begin{aligned} & \sum_{\lambda} (p_{\lambda}^{\rho}(\xi)^* p_{\lambda}^{\rho}(\eta) + p_{\lambda}^{\rho}(\eta)^* p_{\lambda}^{\rho}(\xi)) = 2\langle \xi, \eta \rangle, \\ (4.13) \quad & \sum_{\lambda} \left\{ \sum_{p=0}^{2q-1} (-\hat{w}(\rho; \lambda))^p \pi_{\rho}(\hat{e}_{2q-1-p}) \right\} (p_{\lambda}^{\rho}(\xi)^* p_{\lambda}^{\rho}(\eta) + p_{\lambda}^{\rho}(\eta)^* p_{\lambda}^{\rho}(\xi)) = 0, \end{aligned}$$

and, for $n = 2m$,

$$(4.14) \quad \sum_{\lambda} \pi_{\lambda}(\text{pf})(p_{\lambda}^{\rho}(\xi)^* p_{\lambda}^{\rho}(\eta) + p_{\lambda}^{\rho}(\eta)^* p_{\lambda}^{\rho}(\xi)) = 2\pi_{\rho}(\text{pf})\langle \xi, \eta \rangle,$$

$$(4.15) \quad \sum_{\lambda} \pi_{\lambda}(\text{pf})(w(\rho; \lambda) - 1)(p_{\lambda}^{\rho}(\xi)^* p_{\lambda}^{\rho}(\eta) + p_{\lambda}^{\rho}(\eta)^* p_{\lambda}^{\rho}(\xi)) = 0.$$

In the rest of this section, we calculate the eigenvalues of c_q on irreducible $\mathfrak{so}(n)$ -modules. Our method is based on [9] and [21]. It is known that the usual Clifford multiplication satisfies $(\pi_{\mu_1}(g)\xi) \cdot = \pi_{\Delta}(g)\xi \cdot \pi_{\Delta}(g^{-1})$ for g in $\text{Spin}(n)$ and ξ in \mathbb{R}^n . The Clifford homomorphism is also compatible with the action of $\text{SO}(n)$ or $\text{Spin}(n)$.

Lemma 4.12. *For g in $\text{SO}(n)$ or $\text{Spin}(n)$ and ξ in \mathbb{C}^n , we have*

$$(4.16) \quad p_{\lambda}^{\rho}(\pi_{\mu_1}(g)\xi) = \pi_{\lambda}(g)p_{\lambda}^{\rho}(\xi)\pi_{\rho}(g^{-1}).$$

Hence, for e_{ij} in $\mathfrak{so}(n)$,

$$(4.17) \quad p_{\lambda}^{\rho}(\pi_{\mu_1}(e_{ij})\xi) = \pi_{\lambda}(e_{ij})p_{\lambda}^{\rho}(\xi) - p_{\lambda}^{\rho}(\xi)\pi_{\rho}(e_{ij}).$$

Proof. Consider the action of g on $V_{\rho} \otimes \mathbb{C}^n = \bigoplus_{\lambda} V_{\lambda}$, and we have

$$\sum_{\lambda} \pi_{\lambda}(g)p_{\lambda}^{\rho}(\xi)\phi = \pi_{\rho} \otimes \pi_{\mu_1}(g)(\phi \otimes \xi) = \pi_{\rho}(g) \otimes \pi_{\mu_1}(g)\xi = \sum_{\lambda} p_{\lambda}^{\rho}(\pi_{\mu_1}(g)\xi)\pi_{\rho}(g)\phi$$

for $\phi \otimes \xi$ in $V_{\rho} \otimes \mathbb{C}^n$. Then we conclude that $p_{\lambda}^{\rho}(\pi_{\mu_1}(g)\xi) = \pi_{\lambda}(g)p_{\lambda}^{\rho}(\xi)\pi_{\rho}(g^{-1})$. \square

The Clifford homomorphism p_{λ}^{ρ} is defined through the projection $\Pi_{\lambda}^{\rho} : V_{\rho} \otimes \mathbb{C}^n \rightarrow V_{\lambda}$. Therefore Π_{λ}^{ρ} is realized with the Clifford homomorphism.

Lemma 4.13. *The orthogonal projection $\Pi_{\lambda}^{\rho} : V_{\rho} \otimes \mathbb{C}^n \rightarrow V_{\lambda} \subset V_{\rho} \otimes \mathbb{C}^n$ is realized as follows:*

$$(4.18) \quad \Pi_{\lambda}^{\rho}(\phi \otimes \xi) = \sum_i p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(\xi)\phi \otimes e_i.$$

Proof. It follows from (4.16) that the following mapping is an $\mathfrak{so}(n)$ -equivariant injection,

$$V_{\lambda} \ni \psi \mapsto \sum_i p_{\lambda}^{\rho}(e_i)^* \psi \otimes e_i \in V_{\rho} \otimes \mathbb{C}^n.$$

Taking into account (4.3), we decompose $\phi \otimes \xi$ as

$$\phi \otimes \xi = \sum_i \langle \xi, e_i \rangle \phi \otimes e_i = \sum_i \sum_{\lambda} p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(\xi)\phi \otimes e_i = \sum_{\lambda} \sum_i p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(\xi)\phi \otimes e_i.$$

Since $\sum_i p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(\xi)\phi \otimes e_i$ is in V_{λ} for each λ , we have the projection formula (4.18). \square

Lemma 4.12 implies that $\sum_i p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(e_i)$ is invariant under the action of $\mathfrak{so}(n)$ and constant on V_{ρ} .

Proposition 4.14. *We set $d(\rho) := \dim V_{\rho}$. Then*

$$\sum_i p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(e_i) = d(\lambda)/d(\rho).$$

The eigenvalues of c_q and \hat{c}_q on irreducible $\mathfrak{so}(n)$ -module V_ρ are

$$(4.19) \quad \pi_\rho(c_q) = \frac{1}{d(\rho)} \sum_\lambda w(\rho; \lambda)^q d(\lambda), \quad \pi_\rho(\hat{c}_q) = \frac{1}{d(\rho)} \sum_\lambda \hat{w}(\rho; \lambda)^q d(\lambda).$$

Moreover, we have a relation for eigenvalues of the Pfaffian element,

$$(4.20) \quad 2m\pi_\rho(\text{pf}) = \frac{1}{d(\rho)} \sum_\lambda \pi_\lambda(\text{pf}) d(\lambda) = \frac{1}{d(\rho)} \sum_\lambda \pi_\lambda(\text{pf}) w(\rho; \lambda) d(\lambda).$$

Proof. Let $\{\phi_\alpha\}_{\alpha=1}^{\dim V_\rho}$ be an orthonormal basis of V_ρ . Taking the trace of Π_λ^ρ , we have

$$\begin{aligned} d(\lambda) &= \sum_{\alpha, i} \langle \Pi_\lambda^\rho(\phi_\alpha \otimes e_i), \phi_\alpha \otimes e_i \rangle = \sum_{\alpha, i, j} \langle p_\lambda^\rho(e_j)^* p_\lambda^\rho(e_i)(\phi_\alpha) \otimes e_j, \phi_\alpha \otimes e_i \rangle \\ &= \sum_{\alpha, i, j} \langle p_\lambda^\rho(e_j)^* p_\lambda^\rho(e_i)(\phi_\alpha), \phi_\alpha \rangle \delta_{ij} = \sum_\alpha \langle \phi_\alpha, \phi_\alpha \rangle \sum_i p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_i) \\ &= d(\rho) \sum_i p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_i). \end{aligned}$$

Thus we obtain $\sum_i p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_i) = d(\lambda)/d(\rho)$ and easily show (4.19) and (4.20). \square

Calculating $d(\rho)$ by Weyl's dimension formula, we have an explicit formula for $\pi_\rho(c_q)$ in [20] and [21]. Though it is not easy to compute $d(\lambda)$ and $d(\rho)$ separately, D. Calderbank, P. Gauduchon and M. Herzlich show that the relative dimension $d(\lambda)/d(\rho)$ can be calculated more easily [9]. We denote by $\{\lambda_i\}_{i=1}^N$ the highest weights in $V_\rho \otimes \mathbb{C}^n$. From Proposition 4.1, we show that N is even only in the following case: (1) n is even and $\rho^m \neq 0$, (2) n is odd and $\rho^m = 1/2$, (3) the exceptional case. By a straightforward calculation, we have

Proposition 4.15 ([9, Lemma 4.5]).

(1) If N is odd,

$$\frac{d(\lambda_j)}{d(\rho)} = (2\hat{w}(\rho; \lambda_j) + 1) \prod_{k \neq j} \frac{\hat{w}(\rho; \lambda_j) + \hat{w}(\rho; \lambda_k)}{\hat{w}(\rho; \lambda_j) - \hat{w}(\rho; \lambda_k)}.$$

(2) If N is even,

$$\frac{d(\lambda_j)}{d(\rho)} = (2\hat{w}(\rho; \lambda_j) - 1) \prod_{k \neq j} \frac{\hat{w}(\rho; \lambda_j) + \hat{w}(\rho; \lambda_k)}{\hat{w}(\rho; \lambda_j) - \hat{w}(\rho; \lambda_k)}.$$

This proposition gives an interesting relation between the conformal weights and the relative dimensions, which is used in Section 7.

Corollary 4.16.

$$(4.21) \quad \sum_j \frac{1}{2\hat{w}(\rho; \lambda_j) - (-1)^N} \frac{d(\lambda_j)}{d(\rho)} = \begin{cases} 1 & \text{if } N \text{ is odd,} \\ 0 & \text{if } N \text{ is even.} \end{cases}$$

5. GRADIENTS ON RIEMANNIAN MANIFOLDS

In this section, we define gradients and study their fundamental properties. We consider only gradients on Riemannian manifolds. The spin case is left to the readers, where the spin connection is used instead of the Levi-Civita connection [11], [19].

Let (M, g) be an n -dimensional oriented Riemannian manifold and let $\mathbf{SO}(M)$ be the principal $\mathbf{SO}(n)$ bundle of the oriented orthonormal frames on M . For an irreducible unitary representation (π_ρ, V_ρ) of $\mathbf{SO}(n)$, we have an associated Hermitian vector bundle $\mathbf{S}_\rho := \mathbf{SO}(M) \times_{\pi_\rho} V_\rho$. The Levi-Civita connection on $\mathbf{SO}(M)$ gives a covariant derivative ∇ on \mathbf{S}_ρ compatible with fiber metric as follows. Let $e = (e_1, \dots, e_n)$ be a local section of $\mathbf{SO}(M)$. For a unitary basis $\{\phi_\alpha\}_\alpha$ of V_ρ , we have a local frame $\{[e, \phi_\alpha]\}_\alpha$ of \mathbf{S}_ρ . With respect to this local trivialization, the covariant derivative ∇ is defined to be

$$(5.1) \quad \nabla := d + \frac{1}{2} \sum_{ij} g(\nabla^T e_i, e_j) \pi_\rho(e_{ij}),$$

where ∇^T is the Levi-Civita connection on the tangent bundle $T(M)$. Since the connection 1-form is skew Hermitian, the derivative ∇ is compatible with fiber metric, that is, $X\langle\phi, \psi\rangle = \langle\nabla_X \phi, \psi\rangle + \langle\phi, \nabla_X \psi\rangle$ for every vector field X .

We shall extend Clifford homomorphisms to bundle homomorphisms. We consider the tensor bundle $\mathbf{S}_\rho \otimes T_{\mathbb{C}}(M) = \mathbf{S}_\rho \otimes (T(M) \otimes \mathbb{C})$ and decompose it as

$$(5.2) \quad \mathbf{S}_\rho \otimes T_{\mathbb{C}}(M) = \bigoplus_{\lambda} \mathbf{S}_\lambda.$$

For each vector field $X = \sum X^i e_i$, we define a bundle homomorphism $p_\lambda^\rho(X)$ in $\Gamma(M, \text{Hom}(\mathbf{S}_\rho, \mathbf{S}_\lambda))$ by

$$p_\lambda^\rho(X) : \mathbf{S}_\rho \ni [e, \phi] \mapsto \sum_i X^i [e, p_\lambda^\rho(e_i) \phi] \in \mathbf{S}_\lambda.$$

From (4.16), we know that this bundle homomorphism is well defined. Furthermore, we can show from (4.17) and (5.1) that

$$\nabla_Y(p_\lambda^\rho(X)\phi) = p_\lambda^\rho(\nabla_X^T Y)\phi + p_\lambda^\rho(X)\nabla_Y\phi$$

for ϕ in $\Gamma(M, \mathbf{S}_\rho)$.

We define geometric first order differential operators depending on the Riemannian metric g on each associated bundle.

Definition 5.1. Decomposing ∇ along (5.2), we have the first order differential operator $D_\lambda^\rho = \Pi_\lambda^\rho \circ \nabla$,

$$D_\lambda^\rho : \Gamma(M, \mathbf{S}_\rho) \xrightarrow{\nabla} \Gamma(M, \mathbf{S}_\rho \otimes T_{\mathbb{C}}^*(M)) \xrightarrow{\simeq} \Gamma(M, \mathbf{S}_\rho \otimes T_{\mathbb{C}}(M)) \xrightarrow{\Pi_\lambda^\rho} \Gamma(M, \mathbf{S}_\lambda),$$

for each λ . Here Π_λ^ρ is the orthogonal projection defined fiberwise from $\mathbf{S}_\rho \otimes T_{\mathbb{C}}(M)$ onto \mathbf{S}_λ . We call this first order differential operator D_λ^ρ the *gradient associated to ρ and λ* .

Example 5.2. Let \mathbf{S}_Δ be the spinor bundle, where $n = 2m + 1$ and $\Delta = ((1/2)_m)$. We have the irreducible decomposition $\mathbf{S}_\Delta \otimes T_{\mathbb{C}}(M) = \mathbf{S}_T \oplus \mathbf{S}_\Delta$, where $T = (3/2, (1/2)_{m-1})$. Then D_Δ^Δ is the Dirac operator and D_T^Δ is the twistor operator up to a normalization.

Example 5.3. The bundle $\Lambda^p(M)$ of (complex) differential forms is realized as an associated vector bundle,

$$(5.3) \quad \begin{cases} \Lambda_{\pm}^n(M) = \mathbf{S}_{(1_{m-1}, \pm 1)} & \text{for } n = 2m \text{ and } p = m, \\ \Lambda^p(M) \simeq \Lambda^{n-p}(M) = \mathbf{S}_{(1_p)} & \text{otherwise.} \end{cases}$$

On this vector bundle, we have three gradients, the conformal Killing operator C , the exterior derivative d , and the interior derivative d^* up to a normalization.

Because the principal symbol of D_{λ}^{ρ} is the Clifford homomorphism p_{λ}^{ρ} , we have a formula of the gradient D_{λ}^{ρ} ,

$$(5.4) \quad D_{\lambda}^{\rho}(\phi) = \Pi_{\lambda}^{\rho} \left(\sum_i \nabla_{e_i} \phi \otimes e_i^* \right) = \sum_i p_{\lambda}^{\rho}(e_i) \nabla_{e_i} \phi.$$

From a similar discussion to the Dirac operator [19], we show that the formal adjoint operator $(D_{\lambda}^{\rho})^*$ of D_{λ}^{ρ} is

$$(D_{\lambda}^{\rho})^* = - \sum_i p_{\lambda}^{\rho}(e_i)^* \nabla_{e_i}.$$

An important feature of the Dirac operator is conformal covariance. The gradients are also conformally covariant operators. Though this fact has been shown by H. D. Fegan [10], we give an explicit proof. We deform the Riemannian metric g conformally as $g' = \exp(2\sigma)g$ for σ in $C^{\infty}(M)$. We denote the objects associated to g' by adding the symbol “'” to them. The orthonormal frame bundle $\mathbf{SO}(M)$ is isomorphic to $\mathbf{SO}'(M)$,

$$\Phi : \mathbf{SO}(M) \ni e = (e_1, \dots, e_n) \mapsto e' = e^{-\sigma}(e_1, \dots, e_n) \in \mathbf{SO}'(M).$$

Then there is a bundle isometry for each ρ ,

$$\Phi_{\rho} : \mathbf{S}_{\rho} = \mathbf{SO}(M) \times_{\pi_{\rho}} V_{\rho} \ni [e, \phi] \mapsto [e', \phi] \in \mathbf{SO}'(M) \times_{\pi_{\rho}} V_{\rho} = \mathbf{S}'_{\rho}$$

such that $\Phi_{\lambda} \circ p_{\lambda}^{\rho}(X) = e^{-\sigma} p_{\lambda}^{\rho}(X) \circ \Phi_{\rho}$ for each vector X . The Levi-Civita connection ∇^T on $T(M)$ changes as

$$\nabla^{T'}_X Y = \nabla^T_X Y + (X\sigma)Y + (Y\sigma)X - g(X, Y)\text{grad}(\sigma),$$

where $\text{grad}(\sigma) := \sum (e_i \sigma) e_i$ is the gradient vector field of σ with respect to g . We show from (5.1) that the covariant derivative on \mathbf{S}_{ρ} changes as

$$\nabla'_X \Phi_{\rho}(\phi) - \Phi_{\rho}(\nabla_X \phi) = \Phi_{\rho} \left(\frac{1}{2} \sum_{ij} ((e_i \sigma)g(X, e_j) - (e_j \sigma)g(X, e_i)) \pi_{\rho}(e_{ij}) \phi \right)$$

for ϕ in $\Gamma(M, \mathbf{S}_{\rho})$. It follows from (4.1) and (5.4) that

$$\begin{aligned} (5.5) \quad & D_{\lambda}^{\rho} \Phi_{\rho}(\phi) \\ &= \Pi_{\lambda}^{\rho} \left(\sum_i \nabla'_{e'_i} \Phi_{\rho}(\phi) \otimes (e'_i)^* \right) = \sum_i \Pi_{\lambda}^{\rho}(\nabla'_{e'_i} \Phi_{\rho}(\phi) \otimes (e_i)^*) = \sum_i p_{\lambda}^{\rho}(e_i) \nabla'_{e_i} \Phi_{\rho}(\phi) \\ &= e^{-\sigma} \Phi_{\lambda} \left(D_{\lambda}^{\rho} \phi + \frac{1}{2} \sum_{ikl} p_{\lambda}^{\rho}(e_i) ((e_k \sigma)g(e_i, e_l) - (e_l \sigma)g(e_i, e_k)) \pi_{\rho}(e_{kl}) \phi \right) \\ &= e^{-\sigma} \Phi_{\lambda} (D_{\lambda}^{\rho} \phi - w(\rho; \lambda) p_{\lambda}^{\rho}(\text{grad}(\sigma)) \phi). \end{aligned}$$

Here, precisely speaking, D_λ^ρ is defined through not $\Pi_\lambda^\rho : \mathbf{S}_\rho \otimes T_\mathbb{C}(M) \rightarrow \mathbf{S}_\lambda$ but $\Pi_\lambda^\rho : \mathbf{S}_\rho \otimes T_\mathbb{C}^*(M) \rightarrow \mathbf{S}_\lambda$. Therefore we use $\Phi_\lambda \circ p_\lambda^\rho(e_i) = e^\sigma p_\lambda^\rho(e_i) \circ \Phi_\rho$ in the above equation. We also have

$$(5.6) \quad [D_\lambda^\rho, f] = D_\lambda^\rho \circ f - f D_\lambda^\rho = p_\lambda^\rho(\text{grad}(f))$$

for f in $C^\infty(M)$. The equations (5.5) and (5.6) give conformal covariance of D_λ^ρ . The next proposition answers why we call $w(\rho; \lambda)$ the conformal weight.

Proposition 5.4 ([10, Theorem 1.1]). *When we change the Riemannian metric g to $g' = e^{2\sigma}g$, the gradient D_λ^ρ changes as*

$$D_\lambda'^\rho = \left(e^{(w(\rho; \lambda) - 1)\sigma} \Phi_\lambda \right) \circ D_\lambda^\rho \circ \left(e^{w(\rho; \lambda)\sigma} \Phi_\rho \right)^{-1}.$$

In particular, if the dimension of $\ker D_\lambda^\rho$ is finite, then $\dim \ker D_\lambda^\rho$ is a conformal invariant of M .

6. CURVATURE ENDOMORPHISMS

Let R_T be the Riemannian curvature on $T(M)$. For a local oriented orthonormal frame $e = (e_1, \dots, e_n)$, we set a local expression of R_T by $R_{ijkl} := g(R_T(e_i, e_j)e_k, e_l)$, and denote the Ricci tensor by $R_{ij} = \sum_k R_{ikkj}$ and the scalar curvature by $\kappa = \sum_i R_{ii}$. We decompose the Riemannian curvature R_{ijkl} ,

$$R_{ijkl} = W_{ijkl} + K_{ijkl} + S_{ijkl},$$

where

$$(6.1) \quad \begin{aligned} S_{ijkl} &:= \frac{\kappa}{n(n-1)}(\delta_{il}\delta_{jk} - \delta_{ik}\delta_{jl}), \\ E_{ij} &:= \frac{1}{n-2} \left(\frac{\kappa}{n}\delta_{ij} - R_{ij} \right), \\ K_{ijkl} &:= E_{ik}\delta_{jl} + E_{jl}\delta_{ik} - E_{il}\delta_{jk} - E_{jk}\delta_{il}, \\ W_{ijkl} &:= R_{ijkl} - E_{ijkl} - S_{ijkl}. \end{aligned}$$

The conformal Weyl tensor W_{ijkl} and the Einstein tensor E_{ij} satisfy

$$(6.2) \quad \sum_i W_{ijil} = 0, \quad E_{ij} = E_{ji}, \quad \sum_i E_{ii} = 0.$$

We shall discuss curvature endomorphisms on the associate vector bundle \mathbf{S}_ρ . We define the second order derivative $\nabla_{X,Y}^2$ on \mathbf{S}_ρ for vector fields X and Y by

$$\nabla_{X,Y}^2 := \nabla_X \nabla_Y - \nabla_{\nabla_X^T Y}.$$

Then the curvature on \mathbf{S}_ρ is $R_\rho(X, Y) = \nabla_{X,Y}^2 - \nabla_{Y,X}^2$. From (5.1), a local expression of R_ρ is

$$R_\rho(e_i, e_j) = \frac{1}{2} \sum_{ij} R_{ijkl} \pi_\rho(e_{kl}).$$

By an easy calculation, we can decompose the curvature R_ρ ,

$$(6.3) \quad R_\rho(e_i, e_j) = \frac{1}{2} \sum_{kl} W_{ijkl} \pi_\rho(e_{kl}) + \sum_k (E_{ik} \pi_\rho(e_{kj}) - E_{jk} \pi_\rho(e_{ki})) - \frac{\kappa}{n(n-1)} \pi_\rho(e_{ij}).$$

Definition 6.1. We define a curvature endomorphism in $\Gamma(M, \text{End}(\mathbf{S}_\rho))$ by

$$R_\rho^q := \sum_{ij} \pi_\rho(e_{ij}^q) R_\rho(e_i, e_j)$$

for each q . When n is even,

$$R_\rho^{\text{pf}} := \sum_{ij} \pi_\rho(\text{pf}_{ij}) R_\rho(e_i, e_j).$$

Instead of R_ρ^q , we often use the translated curvature endomorphism

$$\hat{R}_\rho^q := \sum_{ij} \pi_\rho(\hat{e}_{ij}^q) R_\rho(e_i, e_j) = \sum_{0 \leq p \leq q} \binom{q}{p} \left(\frac{n-1}{2} \right)^{q-p} R_\rho^p.$$

Example 6.2. $R_{((1/2)_{m-1}, \pm 1/2)}^1$ is $\kappa/4$, and $R_{(1)}^1/2$ is the Ricci transformation.

By Clifford homomorphisms, the curvature endomorphisms are rewritten as

$$\begin{aligned} R_\rho^q &= \sum_{ij, \lambda} w(\rho; \lambda)^q p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_j) R_\rho(e_i, e_j), \\ R_\rho^{\text{pf}} &= \sum_{ij, \lambda} \pi_\lambda(\text{pf}) p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_j) R_\rho(e_i, e_j). \end{aligned}$$

Proposition 6.3. The curvature endomorphisms R_ρ^q and R_ρ^{pf} are self-adjoint endomorphisms of \mathbf{S}_ρ .

Proof. We consider a curvature endomorphism of \mathbf{S}_ρ ,

$$(6.4) \quad R_\rho(\lambda) := \sum_{ij} p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_j) R_\rho(e_i, e_j) = \sum_{ijkl} R_{ijkl} p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_j) \pi_\rho(e_{kl})$$

for each λ . From (4.17), we have

$$\begin{aligned} &\pi_\rho(e_{kl}) p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_j) - p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_j) \pi_\rho(e_{kl}) \\ &= p_\lambda^\rho(\delta_{ki} e_l - \delta_{li} e_k) p_\lambda^\rho(e_j) + p_\lambda^\rho(e_i)^* p_\lambda^\rho(\delta_{kj} e_l - \delta_{lj} e_k). \end{aligned}$$

Then it is easy to show that

$$R_\rho(\lambda) = \sum_{ijkl} R_{ijkl} p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_j) \pi_\rho(e_{kl}) = \sum_{ijkl} R_{ijkl} \pi_\rho(e_{kl}) p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_j) = R_\rho(\lambda)^*.$$

Since R_ρ^q and R_ρ^{pf} are linear combinations of $\{R_\rho(\lambda)\}_\lambda$ with real coefficients, R_ρ^q and R_ρ^{pf} are self-adjoint endomorphisms. \square

We decompose the curvature endomorphisms along (6.3). From (3.10) and (6.2), the Einstein part of \hat{R}_ρ^q is

$$\sum_{ijk} \pi_\rho(\hat{e}_{ij}^q) (E_{ik} \pi_\rho(e_{kj}) - E_{jk} \pi_\rho(e_{ki})) = - \sum_{ij} E_{ij} \pi_\rho(2\hat{e}_{ij}^{q+1} + \hat{e}_{ij}^q).$$

Thus R_ρ^q and \hat{R}_ρ^q decompose as

$$\begin{aligned} R_\rho^q &= \frac{1}{2} \sum_{ijkl} W_{ijkl} \pi_\rho(e_{ij}^q e_{kl}) - \sum_{ij} E_{ij} \pi_\rho(2e_{ij}^{q+1} + ne_{ij}^q) + \frac{\pi_\rho(c_{q+1})\kappa}{n(n-1)}, \\ \hat{R}_\rho^q &= \frac{1}{2} \sum_{ijkl} W_{ijkl} \pi_\rho(\hat{e}_{ij}^q \hat{e}_{kl}) - \sum_{ij} E_{ij} \pi_\rho(2\hat{e}_{ij}^{q+1} + \hat{e}_{ij}^q) + \frac{\pi_\rho(2\hat{c}_{q+1} - (n-1)\hat{c}_q)\kappa}{n(n-1)}. \end{aligned}$$

Example 6.4. If M is the standard sphere S^n , then R_ρ^q is $\pi_\rho(c_{q+1})$.

We consider R_ρ^{pf} . It follows from (4.1), (4.14), and (4.15) that the Einstein part of R_ρ^{pf} is

$$\sum_{\lambda, ijk} \pi_\lambda(\text{pf}) p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_j) (E_{ik} \pi_\rho(e_{kj}) - E_{jk} \pi_\rho(e_{ki})) = 0,$$

and the scalar curvature part is $\pi_\rho(\text{pf})\kappa/(n-1)$.

Proposition 6.5. *The curvature endomorphism R_ρ^{pf} does not depend on the Einstein tensor,*

$$R_\rho^{\text{pf}} = \frac{1}{2} \sum_{ijkl} W_{ijkl} \pi_\rho(\text{pf}_{ij} e_{kl}) + \frac{\pi_\rho(\text{pf})\kappa}{n-1}.$$

Corollary 6.6. (1) *Suppose that M is an even-dimensional conformally flat manifold; then*

$$R_\rho^{\text{pf}} = \frac{\pi_\rho(\text{pf})\kappa}{n-1}.$$

(2) *For the exceptional case that $\rho^{m-1} > 0$ and $\rho^m = 0$,*

$$R_\rho^{\text{pf}} = \frac{1}{2} \sum_{ijkl} W_{ijkl} \pi_\rho(\text{pf}_{ij} e_{kl}).$$

7. BOCHNER-WEITZENBÖCK FORMULAS

The second order differential operator $(D_\lambda^\rho)^* D_\lambda^\rho$ on \mathbf{S}_ρ is realized as

$$(D_\lambda^\rho)^* D_\lambda^\rho = - \sum_{i,j} p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_j) \nabla_{e_i, e_j}^2.$$

From (4.10), we have

$$\sum_\lambda (D_\lambda^\rho)^* D_\lambda^\rho = - \sum_{\lambda, i, j} p_\lambda^\rho(e_i)^* p_\lambda^\rho(e_j) \nabla_{e_i, e_j}^2 = - \sum_{i, j} \delta_{ij} \nabla_{e_i, e_j}^2 = \nabla^* \nabla,$$

where $\nabla^* \nabla$ is the connection Laplacian on \mathbf{S}_ρ defined by $-\sum_i \nabla_{e_i, e_i}^2$.

The universal Bochner-Weitzenböck formula (3.12) gives

$$\begin{aligned} \hat{R}_\rho^{2q} &= \sum_{i,j} \pi_\rho(\hat{e}_{ij}^{2q}) \left(\nabla_{e_i, e_j}^2 - \nabla_{e_j, e_i}^2 \right) \\ &= - \sum_\lambda \hat{w}(\rho; \lambda)^{2q} (D_\lambda^\rho)^* D_\lambda^\rho - \sum_{i,j} \pi_\rho \left(\hat{e}_{ji}^{2q} + \sum_{p=0}^{2q-1} (-1)^p \hat{c}_{2q-1-p} \hat{e}_{ji}^p \right) \nabla_{e_j, e_i}^2 \\ &= \sum_\lambda \left\{ \sum_{p=0}^{2q-1} \pi_\rho(\hat{c}_{2q-1-p}) (-\hat{w}(\rho; \lambda))^p \right\} (D_\lambda^\rho)^* D_\lambda^\rho. \end{aligned}$$

Similarly (4.14) gives

$$R_\rho^{\text{pf}} = \sum_\lambda 2(\pi_\rho(\text{pf}) - \pi_\lambda(\text{pf})) (D_\lambda^\rho)^* D_\lambda^\rho.$$

We are now in a position to state the main theorem.

Theorem 7.1 (Bochner-Weitzenböck formulas). *Let $\{D_\lambda^\rho\}_\lambda$ be the gradients on \mathbf{S}_ρ , and let $\{(D_\lambda^\rho)^*\}_\lambda$ be their formal adjoints. Then*

$$(7.1) \quad \sum_\lambda (D_\lambda^\rho)^* D_\lambda^\rho = \nabla^* \nabla,$$

$$(7.2) \quad \sum_\lambda \left\{ \sum_{p=0}^{2q-1} \pi_\rho(\hat{c}_{2q-1-p})(-\hat{w}(\rho; \lambda))^p \right\} (D_\lambda^\rho)^* D_\lambda^\rho = \hat{R}_\rho^{2q}, \quad q = 1, 2, \dots.$$

When n is even, we also have

$$(7.3) \quad \sum_\lambda 2(\pi_\rho(\text{pf}) - \pi_\lambda(\text{pf}))(D_\lambda^\rho)^* D_\lambda^\rho = R_\rho^{\text{pf}}.$$

For the exceptional case, setting $\lambda_\pm := \rho \pm \mu_m$, we have

$$(7.4) \quad (D_{\lambda_+}^\rho)^* D_{\lambda_+}^\rho - (D_{\lambda_-}^\rho)^* D_{\lambda_-}^\rho = -\frac{1}{4\pi_{\lambda_+}(\text{pf})} \sum_{ijkl} W_{ijkl} \pi_\rho(\text{pf}_{ij} e_{kl}).$$

Remark 7.2. By using (3.13), we obtain other identities

$$\hat{R}_\rho^{2q+1} = -\sum_\lambda \left\{ 2\hat{w}(\rho; \lambda)^{2q+1} + \hat{w}(\rho; \lambda)^{2q} - \sum_{p=0}^{2q} \pi_\rho(\hat{c}_{2q-p})(-\hat{w}(\rho; \lambda))^p \right\} (D_\lambda^\rho)^* D_\lambda^\rho$$

for $q = 0, 1, \dots$. But, from the discussion below, these formulas are linear dependent on (7.2). The above equation for $q = 0$ gives a formula by P. Gauduchon [12, Appendix B],

$$(7.5) \quad \begin{aligned} -\frac{1}{2} R_\rho^1 &= -\frac{\pi_\rho(c_2)\kappa}{2n(n-1)} + \sum E_{ik} \pi_\rho(e_{ik}^2) - \frac{1}{4} \sum W_{ijkl} \pi_\rho(e_{ij} e_{kl}) \\ &= \sum_\lambda w(\rho; \lambda) (D_\lambda^\rho)^* D_\lambda^\rho. \end{aligned}$$

Note that this formula can be proved from (4.11).

We shall discuss linear independence of Bochner-Weitzenböck formulas (7.2) and (7.3). We assume that there are N gradients $\{D_{\lambda_i}^\rho\}_{i=1}^N$ on \mathbf{S}_ρ . In [3], by using the spectral resolution on the standard sphere, T. Branson shows that there are just $[N/2]$ independent identities such that $\sum_i b_{\lambda_i} (D_{\lambda_i}^\rho)^* D_{\lambda_i}^\rho$ is a curvature endomorphism. Therefore our task is to prove that (7.2) and (7.3) give $[N/2]$ independent identities. Define a vector $v(q)$ consisting of the coefficients in (7.2) by

$$v(q) := \left(\sum_{p=0}^{2q-1} (-1)^p \pi_\rho(\hat{c}_{2q-1-p}) \hat{w}(\rho; \lambda_1)^p, \dots, \sum_{p=0}^{2q-1} (-1)^p \pi_\rho(\hat{c}_{2q-1-p}) \hat{w}(\rho; \lambda_N)^p \right).$$

We decompose $(v(1), v(2), \dots, v(q))$ into the product of a $q \times 2q$ matrix $C(q)$ and a $2q \times N$ matrix $W(q)$ given by

$$C(q) := \begin{pmatrix} \pi_\rho(\hat{c}_1) & -\pi_\rho(\hat{c}_0) & 0 & 0 & \cdots & 0 & 0 \\ \pi_\rho(\hat{c}_3) & -\pi_\rho(\hat{c}_2) & \pi_\rho(\hat{c}_1) & -\pi_\rho(\hat{c}_0) & \cdots & 0 & 0 \\ \cdots & \cdots & \cdots & \cdots & \cdots & \cdots & \cdots \\ \pi_\rho(\hat{c}_{2q-1}) & -\pi_\rho(\hat{c}_{2q-2}) & \cdots & \cdots & \cdots & \pi_\rho(\hat{c}_1) & -\pi_\rho(\hat{c}_0) \end{pmatrix},$$

$$W(q) := \begin{pmatrix} 1 & 1 & \cdots & 1 \\ \hat{w}(\rho; \lambda_1) & \hat{w}(\rho; \lambda_2) & \cdots & \hat{w}(\rho; \lambda_N) \\ \hat{w}(\rho; \lambda_1)^2 & \hat{w}(\rho; \lambda_2)^2 & \cdots & \hat{w}(\rho; \lambda_N)^2 \\ \cdots & \cdots & \cdots & \cdots \\ \hat{w}(\rho; \lambda_1)^{2q-1} & \hat{w}(\rho; \lambda_2)^{2q-1} & \cdots & \hat{w}(\rho; \lambda_N)^{2q-1} \end{pmatrix}.$$

Since the conformal weights are different from each other, the rank of the matrix $(v(1), v(2), \dots, v([N/2])) = C([N/2])W([N/2])$ is $[N/2]$ unless it is the exceptional case. For the exceptional case, the rank of $C([N/2])W([N/2])$ is $[N/2] - 1$. But, there is another formula (7.4) independent of (7.2). Thus we have $[N/2]$ independent Bochner-Weitzenböck formulas.

Corollary 7.3. *The formulas (7.2) and (7.3) give all Bochner-Weitzenböck formulas for gradients.*

In the rest of this section, we discuss a relation between our Bochner-Weitzenböck formulas and Branson's conformally covariant operators. In [4, Theorem 1.1], Branson classifies the second order conformally covariant operators on associated vector bundles. We consider differential operators,

$$(7.6) \quad (D_{\lambda_\pm}^\rho)^* D_{\lambda_\pm}^\rho \quad \text{for the exceptional case,}$$

$$(7.7) \quad \sum_\lambda \frac{1}{w(\rho; \lambda) + \frac{n-2}{2}} (D_\lambda^\rho)^* D_\lambda^\rho + \frac{\kappa}{2(n-1)} \quad \text{otherwise.}$$

In the same manner as in Proposition 5.4 we can show that the above operators are conformally covariant. His important observation is that, when N is even, the operator (7.7) has zeroth order. This fact is equivalent to the next proposition.

Proposition 7.4. *Assume that $n = 2m$ and $\rho^m \neq 0$, or $n = 2m + 1$ and $\rho^m = 1/2$. Then, at the symbol level, we have*

$$(7.8) \quad \sum_\lambda \frac{1}{w(\rho; \lambda) + \frac{n-2}{2}} (p_\lambda^\rho(\xi)^* p_\lambda^\rho(\eta) + p_\lambda^\rho(\eta)^* p_\lambda^\rho(\xi)) = 0 \quad \text{for } \xi \text{ and } \eta \text{ in } \mathbb{R}^n.$$

Remark 7.5. Corollary 7.3 implies that we can realize (7.8) as a linear combination of (4.13) for $q = 1, \dots, N/2$.

For $n = 2m$, we can prove (7.8) directly. We consider the case of $n = 2m$ and $\rho^m \neq 0$. Since $\pi_\rho(\text{pf})$ is not zero, the equation (7.8) follows from (4.12) and (4.14). Indeed, we can rewrite (7.3) as

$$(7.9) \quad \sum_\lambda \frac{1}{w(\rho; \lambda) + \frac{n-2}{2}} (D_\lambda^\rho)^* D_\lambda^\rho + \frac{\kappa}{2(n-1)} = -\frac{1}{4\pi_\rho(\text{pf})} \sum W_{ijkl} \pi_\rho(\text{pf}_{ij} e_{kl}).$$

Thus the Bochner-Weitzenböck formula associated to the Pfaffian element is just Branson's conformally covariant operator with zeroth order. For the exceptional case, we also have a zeroth order conformally covariant operator, $(D_{\lambda_+}^\rho)^* D_{\lambda_+}^\rho - (D_{\lambda_-}^\rho)^* D_{\lambda_-}^\rho$ in (7.4).

For $n = 2m + 1$, since there is no Casimir element corresponding to the Pfaffian element, the author does not now have a direct proof of (7.8). Nevertheless, it would be significant to prove that the equation (7.8) provides a Bochner-Weitzenböck formula omitting the Einstein tensor. Let $R_\rho(\lambda)$ be the curvature endomorphism given by (6.4). It follows from (7.8) that

$$\begin{aligned} \sum_{\lambda} \frac{1}{w(\rho; \lambda) + \frac{n-2}{2}} (D_{\lambda}^{\rho})^* D_{\lambda}^{\rho} &= -\frac{1}{2} \sum_{\lambda} \frac{1}{w(\rho; \lambda) + \frac{n-2}{2}} p_{\lambda}^{\rho}(e_i) p_{\lambda}^{\rho}(e_j) (\nabla_{e_i, e_j}^2 - \nabla_{e_j, e_i}^2) \\ &= -\frac{1}{2} \sum_{\lambda} \frac{1}{w(\rho; \lambda) + \frac{n-2}{2}} R_{\rho}(\lambda). \end{aligned}$$

Then we shall investigate the curvature endomorphism

$$(7.10) \quad \sum_{\lambda} \frac{1}{w(\rho; \lambda) + \frac{n-2}{2}} R_{\rho}(\lambda).$$

It is easy to see that $R_{\rho}(\lambda)$ is equal to

$$\frac{1}{2} \sum_{ijkl} W_{ijkl} p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(e_j) \pi_{\rho}(e_{kl}) - (2w(\rho; \lambda) + n) \sum_{i,j} p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(e_j) E_{ij} + \frac{w(\rho; \lambda) d(\lambda) \kappa}{n(n-1)d(\rho)},$$

and the Einstein part in (7.10) vanishes because of (6.2) and (7.8). Moreover, we show from (4.21) that the scalar curvature part in (7.10) is $\kappa/(n-1)$. Thus we have

$$\begin{aligned} (7.11) \quad & \sum_{\lambda} \frac{1}{w(\rho; \lambda) + \frac{n-2}{2}} (D_{\lambda}^{\rho})^* D_{\lambda}^{\rho} + \frac{\kappa}{2(n-1)} \\ &= -\frac{1}{4} \sum_{\lambda, i, j, k, l} \frac{1}{w(\rho; \lambda) + \frac{n-2}{2}} W_{ijkl} p_{\lambda}^{\rho}(e_i)^* p_{\lambda}^{\rho}(e_j) \pi_{\rho}(e_{kl}). \end{aligned}$$

Remark 7.6. We know from [4, Theorem 1.1] that, if N is even, up to a constant multiple, there is a unique optimal Bochner-Weitzenböck formula omitting the Einstein tensor: (7.4) for the exceptional case and (7.11) for other cases. If N is odd, there is no such Bochner-Weitzenböck formula.

8. EXAMPLES

We give examples and applications of our Bochner-Weitzenböck formulas. We often simply denote $D_{\lambda_i}^{\rho}$ by D_i in this section.

Example 8.1 (spinors). We discuss gradients on the spinor bundle \mathbf{S}_{Δ} for $n = 2m + 1$. There are two irreducible components in $\mathbf{S}_{\Delta} \otimes T_{\mathbb{C}}(M)$ whose highest weights are $\lambda_1 = (3/2, (1/2)_{m-1})$ and $\lambda_2 = \Delta = ((1/2)_m)$. From (7.1) and (7.5), we have

$$D_1^* D_1 + D_2^* D_2 = \nabla^* \nabla, \quad \frac{1}{2} D_1^* D_1 - \frac{n-1}{2} D_2^* D_2 = -\frac{1}{2} R_{\Delta}^1 = -\frac{\kappa}{8}.$$

Since the Dirac operator D is $\sqrt{n} D_2$ and the twistor operator T is $\sqrt{n/(n-1)} D_1$, we have

$$D^2 = \nabla^* \nabla + \frac{1}{4} \kappa, \quad D^2 = \frac{n}{4(n-1)} \kappa + T^* T.$$

The first equation gives a vanishing theorem. If (M, g) is a compact spin manifold with positive scalar curvature, then the kernel of D is zero. The second equation gives Friedrich's estimate for eigenvalues of D^2 . Since $T^* T$ is a non-negative operator on M , each eigenvalue μ of D^2 satisfies $\mu \geq \frac{n}{4(n-1)} \min_{x \in M} \kappa(x)$ [11].

Example 8.2 (differential forms). We consider the bundle of differential forms $\mathbf{S}_{(1_p)} = \Lambda^p(M) \otimes \mathbb{C} \simeq \Lambda^{n-p}(M) \otimes \mathbb{C}$. We have three irreducible components in $\mathbf{S}_{(1_p)} \otimes T_{\mathbb{C}}(M)$ whose highest weights are

$$\lambda_1 = (2, 1_{p-1}), \quad \lambda_2 = (1_{p+1}), \quad \lambda_3 = (1_{p-1}).$$

From (7.1) and (7.5), we get

$$\begin{aligned} D_1^* D_1 + D_2^* D_2 + D_3^* D_3 &= \nabla^* \nabla, \\ D_1^* D_1 - p D_2^* D_2 - (n-p) D_3^* D_3 &= -\frac{1}{2} R_{(1_p)}^1. \end{aligned}$$

The operators D_1 , D_2 and D_3 are constant multiples of the conformal Killing operator C , the exterior derivative d and the interior derivative d^* , respectively. We normalize $\{D_i\}_{1 \leq i \leq 3}$ and obtain

$$\begin{aligned} (8.1) \quad C^* C + \frac{1}{p+1} d^* d + \frac{1}{n-p+1} d d^* &= \nabla^* \nabla, \\ C^* C - \frac{p}{p+1} d^* d - \frac{n-p}{n-p+1} d d^* &= -\frac{1}{2} R_{(1_p)}^1. \end{aligned}$$

In particular, we get a well-known formula, $d^* d + d d^* = \nabla^* \nabla + R_{(1_p)}^1/2$.

By using (8.1), we can prove eigenvalue estimates of the Laplace operator $d^* d + d d^*$, Lichnerowicz's estimate for functions and Gallot-Meyer's estimate for differential forms [13]. Suppose that (M, g) is a compact Riemannian manifold. For eigenfunction f of $d^* d$ with non-zero eigenvalue μ , we have

$$\begin{aligned} \mu \|df\|^2 &= (d d^* df, df) = \frac{n}{n-1} \left(\left(C^* C + \frac{1}{2} R_{\Lambda^1}^1 \right) df, df \right) \\ &= \frac{n}{n-1} \{ \|C df\|^2 + (Ric(df), df) \} \geq \frac{n}{n-1} (Ric(df), df), \end{aligned}$$

where (ϕ, ψ) denotes $\int_M \langle \phi, \psi \rangle dv$. Accordingly, non-zero eigenvalue of $d^* d$ on $\Lambda^0(M)$ has a lower bound depending on the Ricci curvature.

Suppose that (M, g) is a compact Riemannian manifold of positive curvature. In other words, there exists a constant $r > 0$ such that $R_{ijkl} \geq r(\delta_{il}\delta_{jk} - \delta_{ik}\delta_{jl})$. For ϕ in $\Gamma(M, \Lambda^p(M))$,

$$\begin{aligned} ((d^* d + d d^*) \phi, \phi) &\geq \frac{n-p+1}{n-p} \left(\left(\frac{p}{p+1} d^* d + \frac{n-p}{n-p+1} d d^* \right) \phi, \phi \right) \\ &= \frac{n-p+1}{n-p} \|C \phi\|^2 + \frac{n-p+1}{2(n-p)} (R_{(1_p)}^1 \phi, \phi) \\ &\geq \frac{n-p+1}{2(n-p)} r \pi_{(1_p)}(c_2) \|\phi\|^2 = p(n-p+1)r \|\phi\|^2. \end{aligned}$$

Therefore the eigenvalue μ of $d^* d + d d^*$ on $\Lambda^p(M)$ satisfies $\mu \geq p(n-p+1)r$.

Example 8.3 (The conformal Weyl tensor). Let (M, g) be an n -dimensional Riemannian manifold with $n \geq 5$. The four-dimensional case is discussed in the next section. The conformal Weyl tensor $W = W_{ijkl}$ is a section of \mathbf{S}_ρ with $\rho = (2_2)$. The highest weights of irreducible components in $\mathbf{S}_\rho \otimes T_{\mathbb{C}}(M)$ are

$$\begin{cases} (3, 2), & (2_2), & (2, 1), & & n = 5, \\ (3, 2), & (2_2, 1), & (2_2, -1), & (2, 1), & n = 6, \\ (3, 2) & (2_2, 1), & (2, 1), & & n \geq 7. \end{cases}$$

Because of the second Bianchi identity, the projections of ∇W to $\mathbf{S}_{(2_2, \pm 1)}$ and $\mathbf{S}_{(2_2)}$ are zero [23]. Setting $\lambda_1 := (3, 2)$ and $\lambda_2 := (2, 1)$, we show from (7.1) and (7.5) that there is a Bochner-Weitzenböck formula for the conformal Weyl tensor,

$$\frac{n+1}{2} D_2^* D_2 W = \nabla^* \nabla W + \frac{1}{4} R_\rho^1 W.$$

It is known that $D_2 W$ is a constant multiple of $\sum_s (\nabla^s W_{sijk} + \nabla^s W_{sjik})$ [23]. Then $D_2 W = 0$ is equivalent to $\delta W = 0$, where $\delta W := -\sum \nabla^s W_{sijk}$. If (M, g) is a Riemannian manifold with $\delta W = 0$, then W satisfies

$$\nabla^* \nabla W + \frac{1}{4} R_\rho^1 W = 0.$$

Example 8.4 (The exceptional case). We consider the exceptional case that $n = 2m$ and $\rho^{m-1} > \rho^m = 0$. When ρ is (1_{m-1}) , $D_{\lambda_\pm}^\rho$ is a constant multiple of d_\pm on $\Lambda^{m-1}(M) \otimes \mathbb{C}$, where d_\pm denotes $\frac{1}{2}(1 \pm *)d$ for $m = 2l$ and $\frac{1}{2}(1 \pm \sqrt{-1}*)d$ for $m = 2l + 1$. Then we have $d_+^* d_+ = d_-^* d_-$. But, in general, $(D_{\lambda_+}^\rho)^* D_{\lambda_+}^\rho$ is not equal to $(D_{\lambda_-}^\rho)^* D_{\lambda_-}^\rho$. If (M, g) is a conformally flat manifold, then we have $(D_{\lambda_+}^\rho)^* D_{\lambda_+}^\rho = (D_{\lambda_-}^\rho)^* D_{\lambda_-}^\rho$ from (7.4).

Example 8.5 (Branson-Hijazi vanishing theorem). We consider the case of $n = 2m$ and $\rho = (p_{m-1}, \pm p)$ for $p = 1, 3/2, 2, 5/2, \dots$. We have two irreducible components in $\mathbf{S}_\rho \otimes T_\mathbb{C}(M)$ and set $\lambda_1 := (p+1, p_{m-2}, \pm p)$ and $\lambda_2 := (p_{m-1}, \pm p \mp 1)$. Then we have

$$\begin{aligned} D_1^* D_1 + D_2^* D_2 &= \nabla^* \nabla, \\ p D_1^* D_1 + (-p - m + 1) D_2^* D_2 &= -\frac{1}{2} R_\rho^1, \\ 2(\pi_\rho(\text{pf}) - \pi_{\lambda_1}(\text{pf})) D_1^* D_1 + 2(\pi_\rho(\text{pf}) - \pi_{\lambda_2}(\text{pf})) D_2^* D_2 &= R_\rho^{\text{pf}}. \end{aligned}$$

The second identity is linear dependent on the third one. We actually have $\pi_\rho(\text{pf}) R_\rho^1 = p(p+m-1) R_\rho^{\text{pf}}$. Since R_ρ^{pf} omits the Einstein tensor, we have

$$(8.2) \quad \frac{2p+m-1}{p} D_2^* D_2 = \nabla^* \nabla + \frac{p+m-1}{4m-2} \kappa + \frac{1}{4p} \sum W_{ijkl} \pi_\rho(e_{ij} e_{kl}).$$

T. Branson and O. Hijazi give a vanishing theorem as follows [7, Theorem 4.2]. We diagonalize W_{ijkl} such that $W_{ijkl} = c_{ij}(x)(\delta_{il}\delta_{jk} - \delta_{ik}\delta_{jl})$ on each point x in M . Then

$$\pi_\rho(c_2) \min_{i < j} c_{ij}(x) \leq \frac{1}{2} \sum_{ijkl} W_{ijkl} \pi_\rho(e_{ij} e_{kl}) = \sum_{ij} c_{ij} \pi_\rho(e_{ij} e_{ji}) \leq \pi_\rho(c_2) \max_{i < j} c_{ij}(x).$$

Supposing $\frac{\kappa(x)}{n(n-1)} + \min c_{ij}(x) > 0$ for all x in M , we have $\ker D_2 = 0$ by (8.2). In particular, for $\rho = (1_{m-1}, \pm 1)$, we have Bourguignon's vanishing theorem, $H^m(M, \mathbb{R}) = 0$ in [2]. They also discuss a relation to the first eigenvalue of Yamabe Laplacian in [6], [7].

Through Examples 8.2 and 8.5, we note that it needs to investigate curvature endomorphisms for vanishing theorems and eigenvalue estimates. We diagonalize R_{ijkl} such that $R_{ijkl} = r_{ij}(x)(\delta_{il}\delta_{jk} - \delta_{ik}\delta_{jl})$. Then we have $\hat{R}_\rho^q = \sum_{ij} r_{ij} \pi_\rho(\hat{e}_{ij}^q \hat{e}_{ji})$.

The curvature endomorphism \hat{R}_ρ^q for $q = 1, 2$ can be estimated,

$$\pi_\rho(c_2) \min_{i < j} r_{ij}(x) \leq \hat{R}_\rho^1 = \frac{2}{n} \hat{R}_\rho^2 = \sum_{ij} r_{ij} \pi_\rho(e_{ij} e_{ji}) \leq \pi_\rho(c_2) \max_{i < j} r_{ij}(x).$$

Since we do not know whether $\pi_\rho(\hat{e}_{ij}^q \hat{e}_{ji})$ for $q \geq 3$ is non-negative, it is difficult to estimate \hat{R}_ρ^q for $q \geq 3$. But, in the 4-dimensional case, we can estimate the curvature endomorphisms more precisely using the decomposition $\Lambda^2(M) = \Lambda_+^2(M) \oplus \Lambda_-^2(M)$.

Remark 8.6. It is possible to estimate the Einstein part and the scalar curvature part of curvature endomorphisms. For example, the Einstein part of $R_\rho^3 - nR_\rho^2/2 + n^2R_\rho^1/4$ is $-\sum_{ij} 2\pi_\rho(e_{ij}^4)E_{ij}$, which we can estimate because of $\pi_\rho(e_{ij}^4) \geq 0$.

9. THE FOUR-DIMENSIONAL CASE

In this section, we discuss the four-dimensional case. We begin with basic facts of four-dimensional geometry. The Hodge star operator $*$ decomposes $\Lambda^2(\mathbb{R}^4)$ into the self-dual part Λ_+^2 and the anti-self-dual part Λ_-^2 . For an oriented orthonormal basis $e = (e_1, e_2, e_3, e_4)$ of \mathbb{R}^4 , we set a basis of Λ_\pm^2 by

$$X_1^\pm := \frac{1}{2}(e_{14} \pm e_{23}), \quad X_2^\pm := -\frac{1}{2}(e_{13} \pm e_{42}), \quad X_3^\pm := \frac{1}{2}(e_{12} \pm e_{34}).$$

Identifying $\Lambda^2(\mathbb{R}^4)$ with $\mathfrak{so}(4) = \mathfrak{so}(3) \oplus \mathfrak{so}(3)$, we get

$$[X_i^\pm, X_j^\pm] = \sum_{1 \leq k \leq 3} \epsilon_{ijk} X_k^\pm, \quad [X_i^+, X_j^-] = 0,$$

where

$$\epsilon_{ijk} = \begin{cases} \text{sgn} \begin{pmatrix} 1 & 2 & 3 \\ i & j & k \end{pmatrix}, & \{i, j, k\} = \{1, 2, 3\}, \\ 0, & \text{otherwise.} \end{cases}$$

Thus $\{X_i^\pm\}_i$ is a standard basis of $\mathfrak{so}(3)$,

$$(9.1) \quad X_1^\pm = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix}, \quad X_2^\pm = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ -1 & 0 & 0 \end{pmatrix}, \quad X_3^\pm = \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

Let (M, g) be a four-dimensional oriented Riemannian or spin manifold. We consider the Riemannian curvature R_T as an endomorphism of $\Lambda^2(M)$,

$$R_T : \Lambda^2(M) \ni e_{ij} \mapsto \frac{1}{2} \sum R_{ijkl} e_{kl} \in \Lambda^2(M).$$

Then we realize R_T as a 6×6 matrix with respect to basis $\{X_i^+\}_i \cup \{X_j^-\}_j$,

$$\begin{pmatrix} W^+ & 0 \\ 0 & W^- \end{pmatrix} + \begin{pmatrix} 0 & K \\ K^t & 0 \end{pmatrix} + \begin{pmatrix} -\kappa/12 & 0 \\ 0 & -\kappa/12 \end{pmatrix},$$

where W^+ (resp. W^-) is the self-dual (resp. anti-self-dual) conformal Weyl tensor and K corresponds to K_{ijkl} in (6.1). In other words, we have

$$\begin{aligned} R_T(X_i^+) &= \sum W_{ij}^+ X_j^+ + \sum K_{ji} X_j^- - \frac{\kappa}{12} X_i^+, \\ R_T(X_i^-) &= \sum W_{ij}^- X_j^- + \sum K_{ij} X_j^+ - \frac{\kappa}{12} X_i^-. \end{aligned}$$

From now on, we denote the highest weight ρ of $\mathfrak{so}(4)$ by $\rho = (\frac{k+l}{2}, \frac{k-l}{2})$ for non-negative integers k and l . Note that ρ corresponds to the highest weight $(k/2) \hat{\otimes} (l/2)$ of $\mathfrak{so}(3) \oplus \mathfrak{so}(3)$.

We define a curvature endomorphism on \mathbf{S}_ρ by

$$R_\rho^\pm := 4 \sum_i \pi_\rho(X_i^\pm R_T(X_i^\pm)).$$

Since $4 \sum X_i^+ X_i^+ = -c_2/2 - 2\text{pf}$ and $4 \sum X_i^- X_i^- = -c_2/2 + 2\text{pf}$, we have

$$\begin{aligned} R_\rho^+ &= 4 \sum_{ij} W_{ij}^+ \pi_\rho(X_i^+ X_j^+) + 4 \sum_{ij} K_{ij} \pi_\rho(X_i^- X_j^+) + \frac{k(k+2)}{12} \kappa, \\ R_\rho^- &= 4 \sum_{ij} W_{ij}^- \pi_\rho(X_i^- X_j^-) + 4 \sum_{ij} K_{ij} \pi_\rho(X_i^- X_j^+) + \frac{l(l+2)}{12} \kappa. \end{aligned}$$

Lemma 9.1. (1) *The curvature endomorphisms R_ρ^1 and R_ρ^{pf} satisfy*

$$(9.2) \quad R_\rho^1 = R_\rho^+ + R_\rho^-, \quad R_\rho^{\text{pf}} = R_\rho^+ - R_\rho^-.$$

(2) *If ρ is $(k/2, k/2)$, then R_ρ^- is zero and R_ρ^+ does not depend on the Einstein tensor. If ρ is $(\frac{k+1}{2}, \frac{k-1}{2})$, then R_ρ^- does not depend on W^- .*

Proof. By using (4.5), we can show (9.2) straightforwardly. We shall prove the second claim. When ρ is $(k/2, k/2)$, $\pi_\rho(X_i^-)$ is zero for each i . Then R_ρ^- is zero and R_ρ^+ does not depend on the Einstein tensor. When ρ is $(\frac{k+1}{2}, \frac{k-1}{2})$, $\{\pi_\rho(X_i^-)\}_{1 \leq i \leq 3}$ gives the spin 1/2 representation of $\mathfrak{so}(3) = \mathfrak{su}(2)$. Then we have $\sum_{ij} W_{ij}^- \pi_\rho(X_i^- X_j^-) = 0$ because of $W_{ij}^- = W_{ji}^-$ and $\sum W_{ii}^- = 0$. \square

The irreducible decomposition of $\mathbf{S}_\rho \otimes T_{\mathbb{C}}(M)$ is $\bigoplus_{i=1}^4 \mathbf{S}_{\lambda_i}$, where

$$\lambda_1 = \rho + \mu_1, \quad \lambda_2 = \rho + \mu_2, \quad \lambda_3 = \rho - \mu_2, \quad \lambda_4 = \rho - \mu_1.$$

From (7.1), (7.2) and (7.3), we have all Bochner-Weitzenböck formulas for the four-dimensional case,

$$\begin{aligned} D_1^* D_1 + D_2^* D_2 + D_3^* D_3 + D_4^* D_4 &= \nabla^* \nabla, \\ k D_1^* D_1 + k D_2^* D_2 - (k+2) D_3^* D_3 - (k+2) D_4^* D_4 &= -R_\rho^+, \\ l D_1^* D_1 - (l+2) D_2^* D_2 + l D_3^* D_3 - (l+2) D_4^* D_4 &= -R_\rho^-. \end{aligned}$$

We shall state some vanishing theorems. First, we consider the case that ρ is $(k/2, k/2)$ for a positive integer k . Then we get $D_2 = D_4 = 0$ and

$$\frac{2(k+1)}{k} D_3^* D_3 = \nabla^* \nabla + \frac{4}{k} \sum W_{ij}^+ \pi_\rho(X_i^+ X_j^+) + \frac{k+2}{12} \kappa.$$

Example 9.2. If M is a compact anti-self-dual manifold, then

$$\frac{2(k+1)}{k} D_3^* D_3 = \nabla^* \nabla + \frac{k+2}{12} \kappa.$$

The kernel of D_3 is isomorphic to a cohomology on the twistor space of M [14].

Example 9.3 (self-dual 2-forms). For $\rho = (1, 1)$, the associated vector bundle \mathbf{S}_ρ is $\Lambda_+^2(M) \otimes \mathbb{C}$ and the kernel of D_3 is the space of harmonic self-dual 2-forms. If ω

is in $\ker D_3$, then we have a formula in [2],

$$0 = \langle \nabla^* \nabla \omega, \omega \rangle + 2 \langle W^+(\omega), \omega \rangle + \frac{1}{3} \kappa \langle \omega, \omega \rangle,$$

where $W^+(\omega) = W^+ (\sum \omega_i X_i^+) = \sum \omega_i W_{ij}^+ X_j^+$.

Example 9.4 (the self-dual conformal Weyl tensor). For $\rho = (2, 2)$, the vector bundle \mathbf{S}_ρ is embedded into $\Lambda_+^2(M) \otimes (\Lambda_+^2(M))^* \otimes \mathbb{C}$ and the sections of \mathbf{S}_ρ are realized locally as trace-free symmetric 3×3 matrices. Therefore we have $\langle Z, Z \rangle = \text{tr}(Z\bar{Z})$ and $\pi_\rho(X_i^+)Z = X_i^+Z - ZX_i^+$ for Z in $\Gamma(M, \mathbf{S}_\rho)$. Here we use the matrix realization (9.1) of $\{X_i^+\}_i$. It follows that

$$\begin{aligned} \sum W_{ij}^+ \pi_\rho(X_i^+ X_j^+) Z &= \sum W_{ij}^+ (X_j^+ X_i^+ Z - X_j^+ Z X_i^+ - X_i^+ Z X_j^+ + Z X_i^+ X_j^+) \\ &= W^+ Z + Z W^+ - 2 \sum W_{ij}^+ X_i^+ Z X_j^+. \end{aligned}$$

Choosing W^+ as a section Z of \mathbf{S}_ρ , we have

$$\begin{aligned} \left\langle \sum W_{ij}^+ \pi_\rho(X_i^+ X_j^+) W^+, W^+ \right\rangle &= 2 \text{tr}((W^+)^3) + 4 \text{tr}((W^+)^3) = 6 \text{tr}((W^+)^3) \\ &= 18 \det(W^+). \end{aligned}$$

If $\delta W^+ = -\sum_s \nabla^s W_{sijk}^+ = 0$, then $D_3 W^+$ is zero and W^+ satisfies the following equation in [1, pp. 451–455]:

$$0 = \langle \nabla^* \nabla W^+, W^+ \rangle + 6 \text{tr}((W^+)^3) + \frac{\kappa}{2} \text{tr}((W^+)^2).$$

Next we consider the case that the highest weight ρ is $(\frac{k+1}{2}, \frac{k-1}{2})$ for a positive integer k . Then we have

$$\begin{aligned} \nabla^* \nabla - 1/3 R_\rho^- &= 4/3 (D_1^* D_1 + D_3^* D_3), \\ \nabla^* \nabla + R_\rho^- &= 4 (D_2^* D_2 + D_4^* D_4), \end{aligned}$$

where

$$R_\rho^- = 4 \sum K_{ij} \pi_\rho(X_i^- X_j^+) + \kappa/4.$$

Example 9.5. If (M, g) is a compact Einstein manifold with positive (resp. negative) constant scalar curvature, then $\ker D_2 \cap \ker D_4$ (resp. $\ker D_1 \cap \ker D_3$) is zero.

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